Physics of topological excitations in Bose -Einstein condensates

Michikazu Kobayashi (University of Tokyo)

Collaborators : Yuki Kawaguchi and Masahito Ueda (University of Tokyo)
Muneto Nitta (Keio University)

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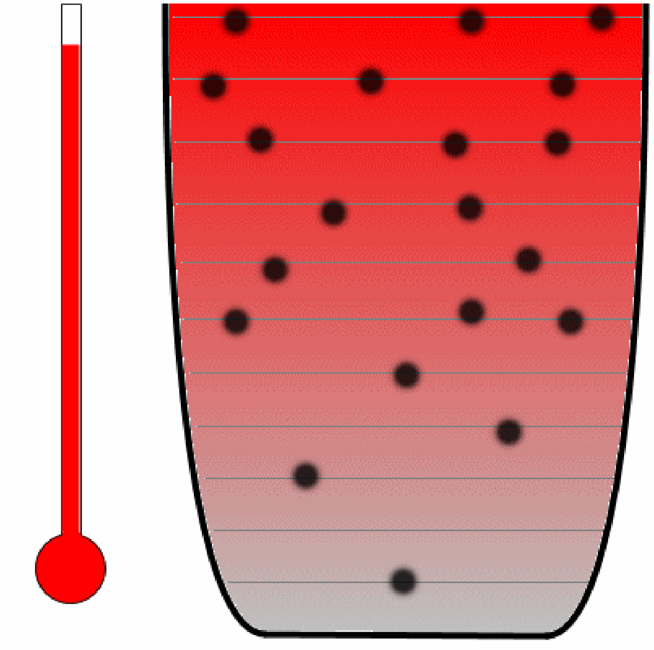
Bose-Einstein condensates



In Bose system at low temperatures, macroscopic numbers of particle occupy the single-particle quantum state, forming the macroscopic wave function.

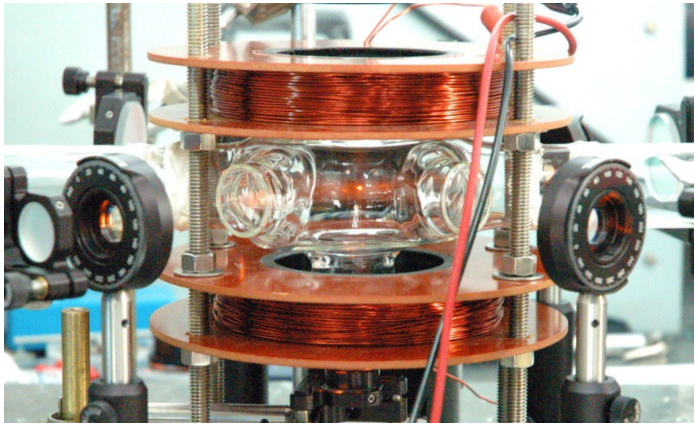
Atoms form the coherent state (origin of superfluidity)

macroscopic wave function : $\langle \hat{\Psi}(x)\hat{\Psi}^\dagger(y)\rangle \xrightarrow{|x-y|\to\infty} \Psi(x)\Psi^*(y)$

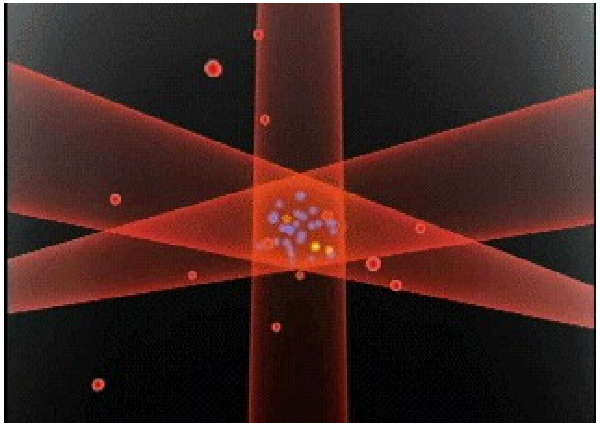
Macroscopic wave function $\Psi(x)$ is the order parameter of BEC

Atomic Bose-Einstein condensates

Dilute alkali atomic Bose-Einstein condensates has been realized in 1997

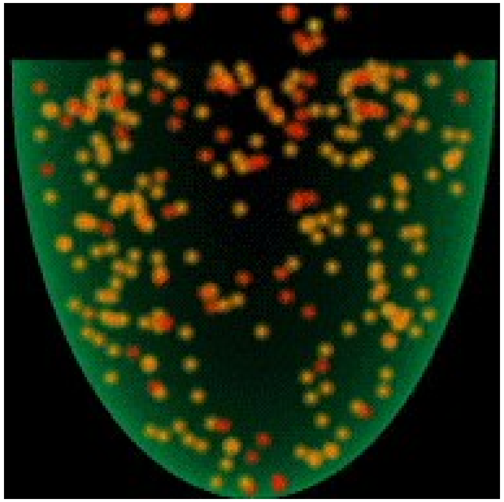


Trap of atoms



Laser cooling of atoms

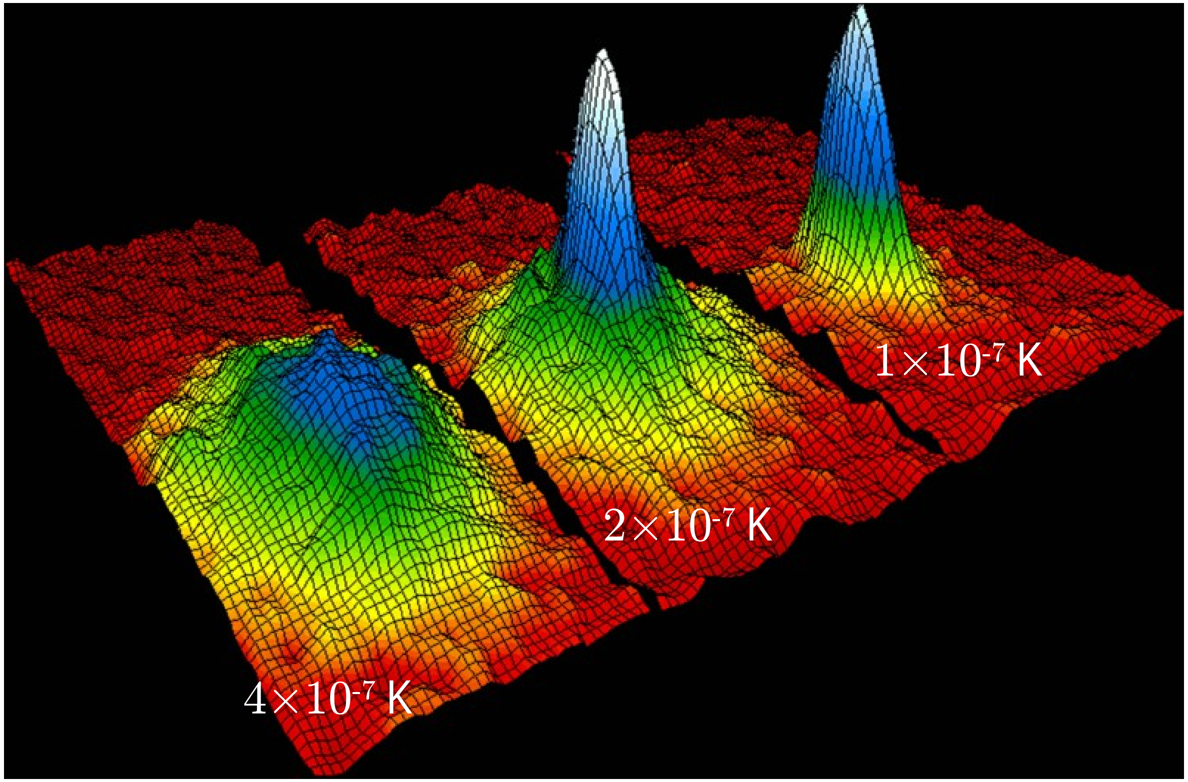
^{87}Rb, ^{23}Na, ^{7}Li, ^{1}H, ^{85}Rb, ^{41}K, ^{4}He, ^{133}Cs, ^{174}Yb, ^{52}Cr, ^{40}Ca, ^{84}Sr



Evaporating cooling of atoms

Atomic Bose-Einstein condensates

BEC of ^{87}Rb



BEC and breaking of gauge symmetry

BEC : appearance of macroscopic wave function (order parameter)

$$\langle \hat{\Psi}(x)\hat{\Psi}^\dagger(y)\rangle \xrightarrow{|x-y|\to\infty} \Psi(x)\Psi^*(y)$$

$\Psi(x) = |\Psi(x)| \exp[i\,\varphi(x)]$: Phase of the complex wave function $\varphi(x)$ is fixed ($U(1)$ gauge symmetry breaking)

Mean-field Hamiltonian at the zero temperature

$$H = \int d\boldsymbol{x} \left[\frac{\hbar^2}{2M} \nabla\Psi^*(\boldsymbol{x})\nabla\Psi(\boldsymbol{x}) + \frac{c_0}{2}|\Psi(\boldsymbol{x})|^4 \right]$$

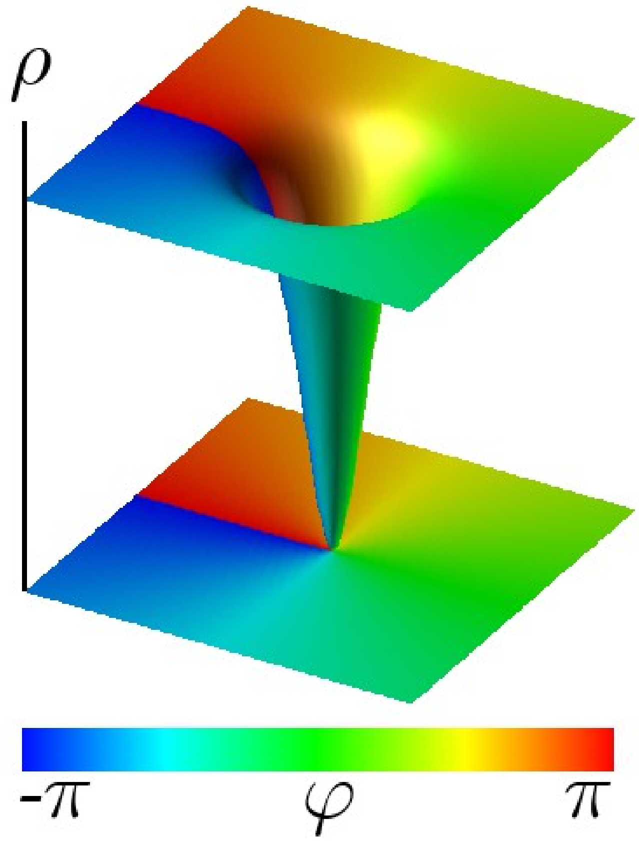
Topological excitations (quantized vortices)

$$\Psi(\boldsymbol{x}) = |\Psi(\boldsymbol{x})| \exp[i\varphi(\boldsymbol{x})]$$

$$\rho(\boldsymbol{x}) = |\Psi(\boldsymbol{x})|^2 \; : \; \text{Fluid density}$$

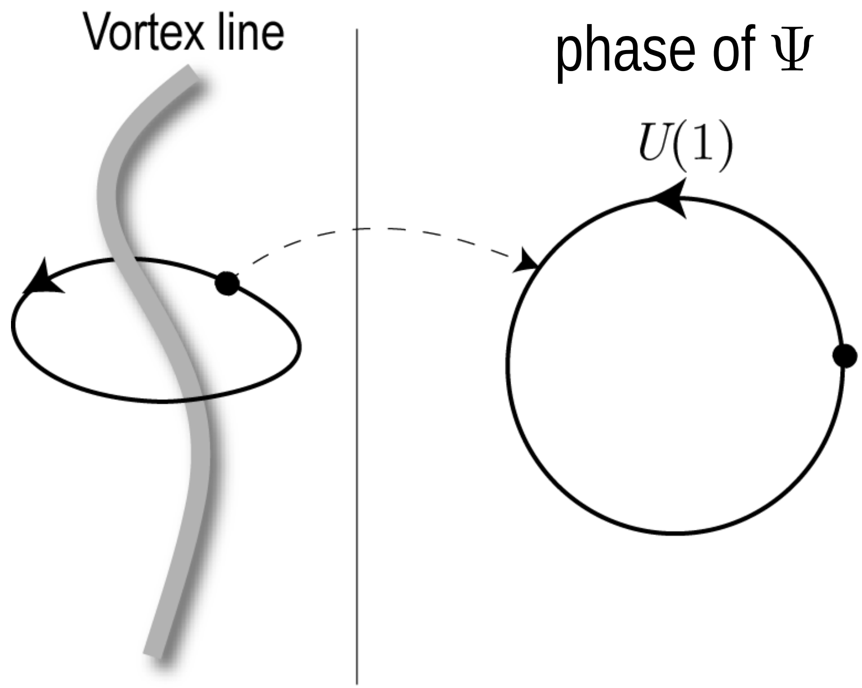
$$\boldsymbol{v}(\boldsymbol{x}) = \frac{\hbar}{m} \nabla \varphi(\boldsymbol{x}) \; : \; \text{Fluid velocity}$$

Topological excitations (defects) of the wave function $(\rho(\boldsymbol{x}) = 0)$ around which phase $\varphi(\boldsymbol{x})$ changes by integer multiple of 2π : **quantized vortices**



Topology of quantized vortices

Topological invariant of vortex is described by fundamental group

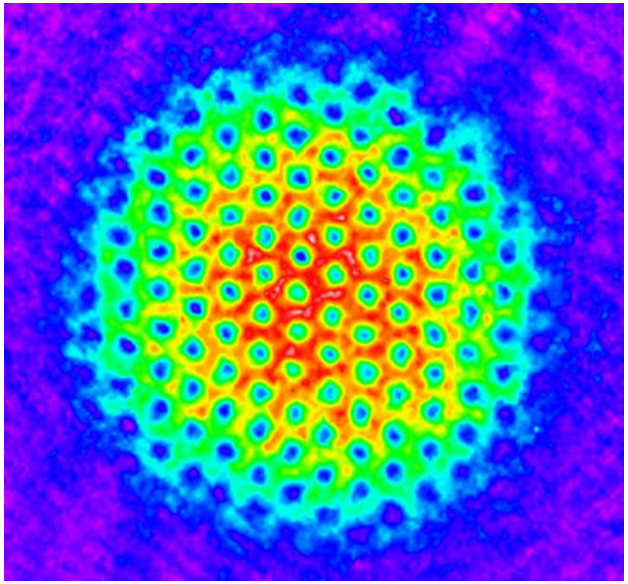


Topological invariants are characterized by how many times Ψ wind the $U(1)$ phase : $\pi_1[U(1)] = \mathbb{Z}$

Experimental observation of vortices

Vortex lattice formation in atomic BEC

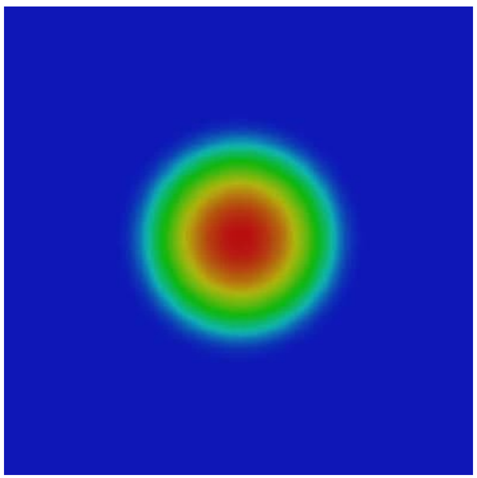
K. W. Madison et al. PRL **86**, 4443 (2001)

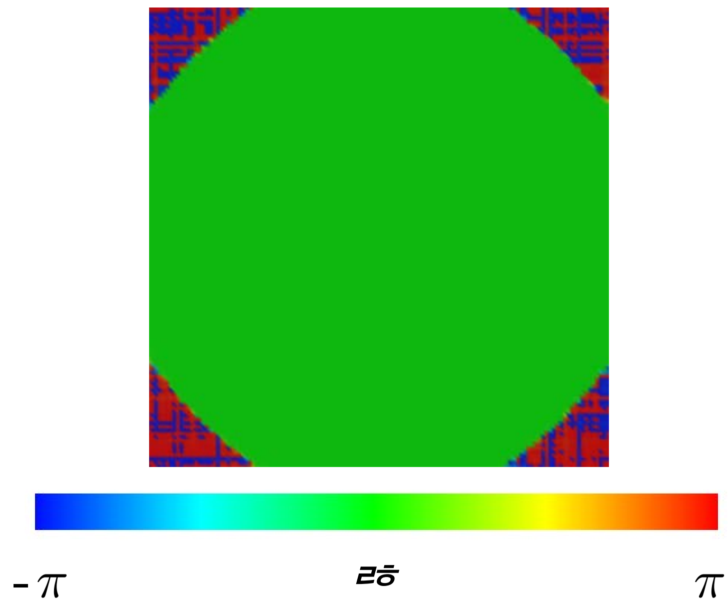


Vortex lattice in ^{87}Rb BEC

Experimental observation of vortices

Vortex lattice formation in atomic BEC



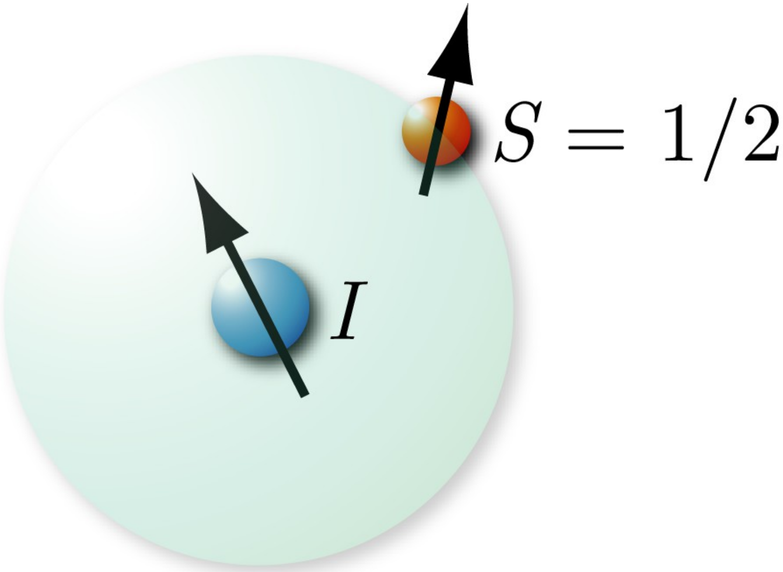


Simulation of vortex lattice formation

Spinor BEC

BEC with spin degrees of freedom

Hyperfine coupling of electron end nuclear spin
$(F = I + L + S)$



^{87}Rb, ^{23}Na, ^{7}Li, ^{41}K	F=1, 2
^{85}Rb	F=2, 3
^{133}Cs	F=3, 4
^{52}Cr	S=3, I=0

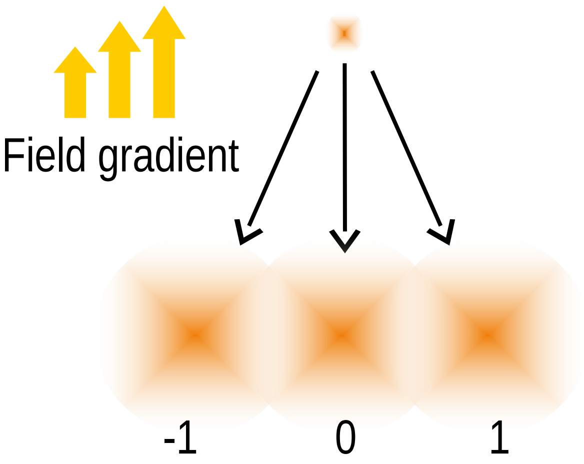
Spinor BEC

BEC with spin degrees of freedom

^{87}Rb $(I = 3/2)$

$$F = 2 \begin{cases} m_F = 2 \\ m_F = 1 \\ m_F = 0 \\ m_F = -1 \\ m_F = -2 \end{cases} \qquad F = 1 \begin{cases} m_F = 1 \\ m_F = 0 \\ m_F = -1 \end{cases}$$

Multicomponent BEC labeled by magnetic sublevel m_F



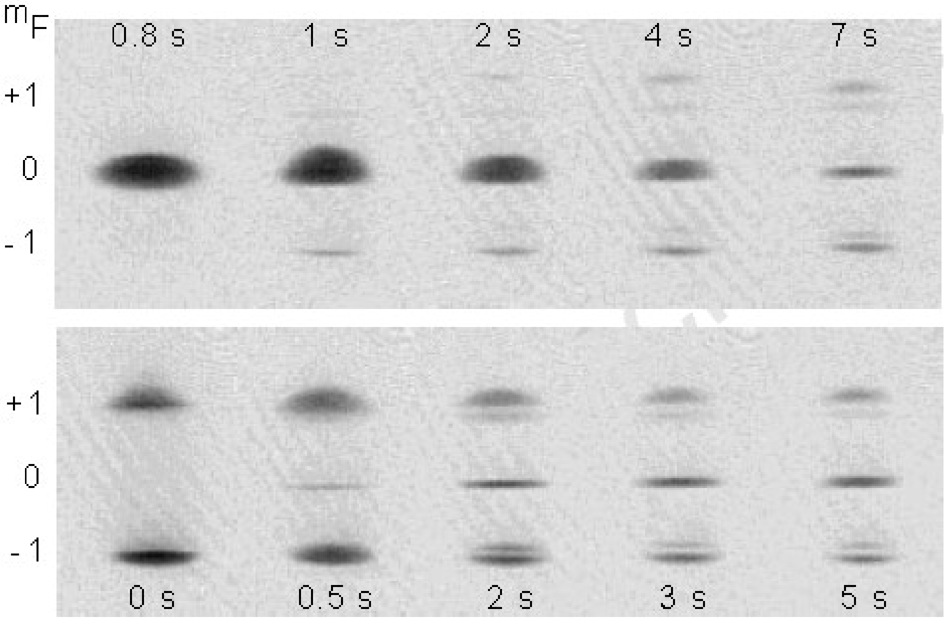
Spin 1 : 3-component BEC

$$\Psi = (\psi_1, \psi_0, \psi_{-1})$$

Spinor BEC

Stern-Gerlach experiment

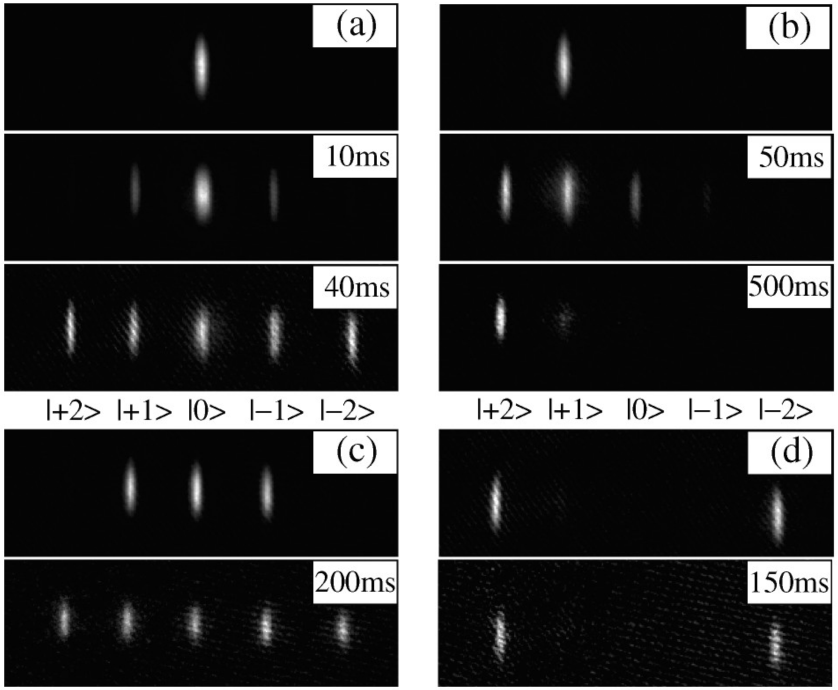
$F = 1$



J. Stenger et al. Nature **396**, 345 (1998)

Rotation of Spin can be observed

$F = 2$



H. Schmaljohann et al. PRL **92**, 040402 (2004)

Theory of Spinor BEC

Hamiltonian of spinor Bosons

$$H = -\int d\boldsymbol{x}\, \frac{\hbar^2}{2M} \nabla\Psi_m^\dagger(\boldsymbol{x})\nabla\Psi_m(\boldsymbol{x})$$
$$+\frac{1}{2}\int d\boldsymbol{x}_1 \int d\boldsymbol{x}_2 \Psi_{m_1}^\dagger(\boldsymbol{x}_1)\Psi_{m_2}^\dagger(\boldsymbol{x}_2) V_{m_1 m_2 m_1' m_x'}(\boldsymbol{x}_1-\boldsymbol{x}_2)\Psi_{m_2'}(\boldsymbol{x}_2)\Psi_{m_1'}(\boldsymbol{x}_1)$$

Low energy contact interaction $(l=0)$

$$V_{m_1 m_2 m_1' m_x'}(\boldsymbol{x}_1-\boldsymbol{x}_2) = \delta(\boldsymbol{x}_1-\boldsymbol{x}_2) \sum_{F=0,2,4} g_F \sum_{m_1,m_2,m_1',m_2',M} O^{F,M}_{m_1 m_2} \left(O^{F,M}_{m_1' m_2'}\right)^*$$

Mean-field Hamiltonian (spin-1)

$$H = \int d\boldsymbol{x} \left[\frac{\hbar^2}{2M} \sum_{m=-1}^{1} \nabla \Psi_m^* \nabla \Psi_m + \frac{c_0}{2}\rho^2 + \frac{c_1}{2}\boldsymbol{F}^2 \right]$$

Density Spin

$$F_z = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix}, \quad F_+ = \begin{pmatrix} 0 & \sqrt{2} & 0 \\ 0 & 0 & \sqrt{2} \\ 0 & 0 & 0 \end{pmatrix}$$

$$F_- = F_+^T, \quad F_x = \frac{F_+ + F_-}{2}, \quad F_y = \frac{F_+ - F_-}{2i}$$

Gauge and spin rotation symmetry of wave function are broken

$$\Psi' = e^{i\varphi}\, e^{-i\boldsymbol{n}\cdot\boldsymbol{F}\alpha}\, \Psi \quad (U(1) \times SO(3))$$

Possible phase

$$H = \int d\boldsymbol{x} \left[\frac{\hbar^2}{2M} \sum_{m=-1}^{1} \nabla \Psi_m^* \nabla \Psi_m + \frac{c_0}{2} \rho^2 + \frac{c_1}{2} \boldsymbol{F}^2 \right]$$

$c_1 > 0$: polar (^{23}Na BEC)

$$e^{i\varphi} e^{-i\boldsymbol{n}\cdot\hat{\boldsymbol{F}}\alpha} \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}$$

$$\boldsymbol{F} = 0$$

$c_1 < 0$: Ferromagnetic (^{87}Rb BEC)

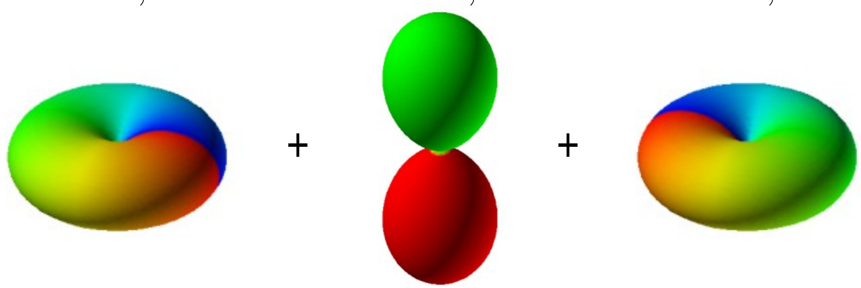
$$e^{i\varphi} e^{-i\boldsymbol{n}\cdot\hat{\boldsymbol{F}}\alpha} \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}$$

$$\boldsymbol{F} \neq 0$$

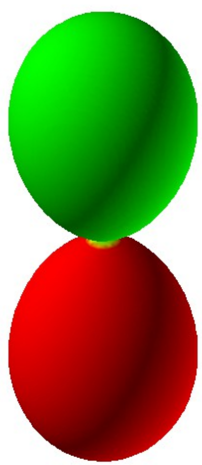
Graphical image by the spherical Harmonics

$$\sum_{m=-1}^{1} \Psi_m Y_{2,m}$$

$Y_{1,1}$ $Y_{1,0}$ $Y_{1,-1}$



$\cos\theta$



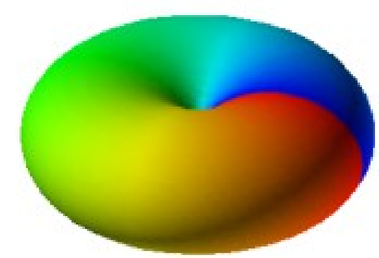
0

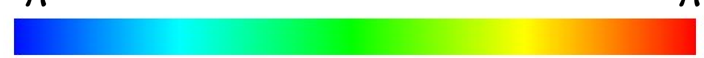
π

Polar

Ferromagnetic

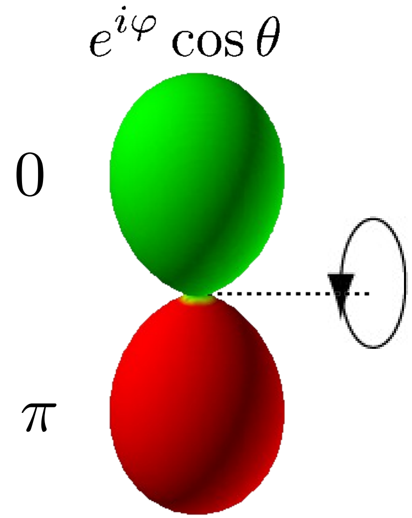
$-e^{i\phi}\sin\theta$

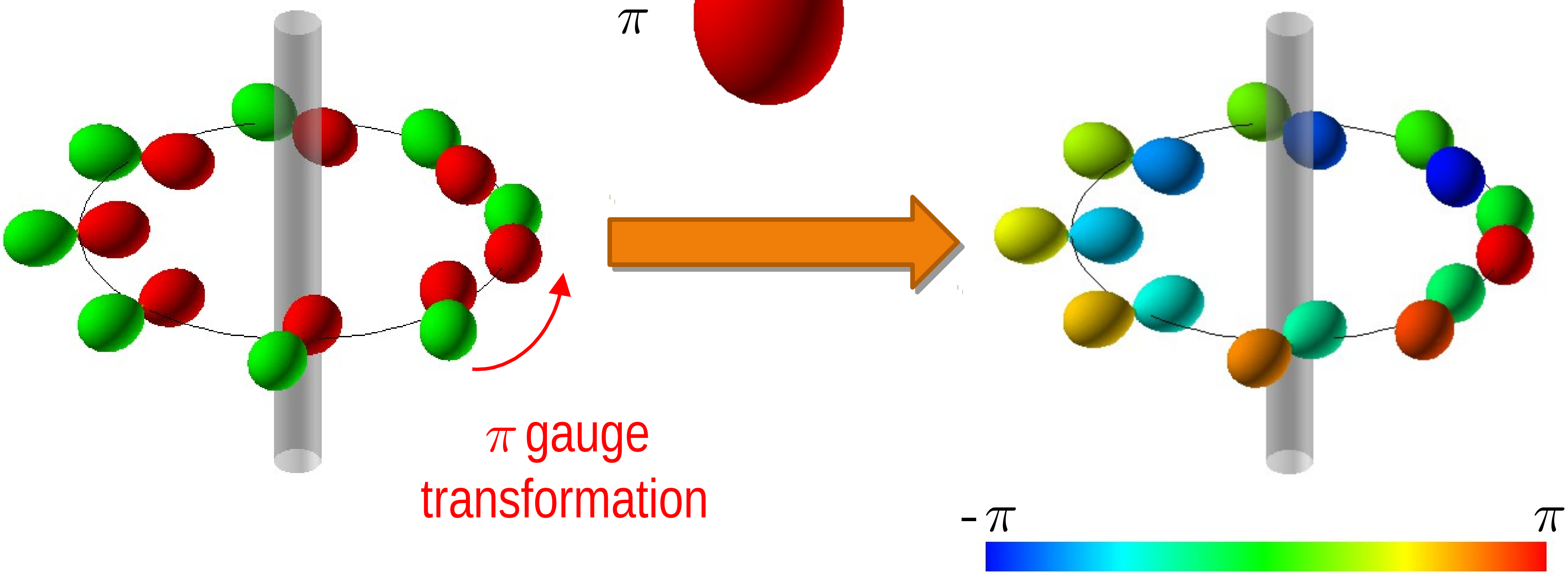




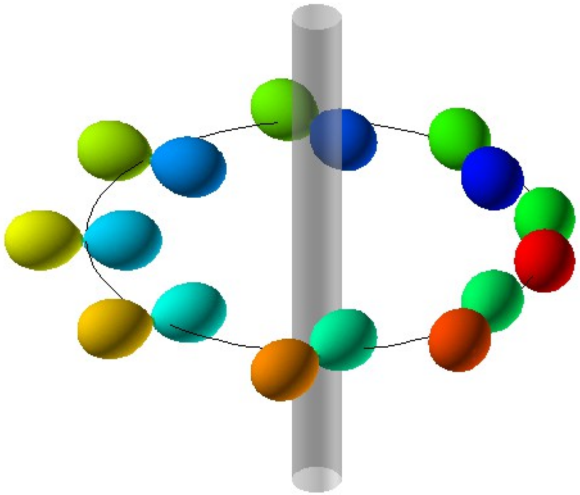
Topological excitation in polar state

$$\frac{G}{H} \simeq \frac{U(1)_\varphi \times S_F^2}{(\mathbb{Z}_2)_{\varphi+F}}$$



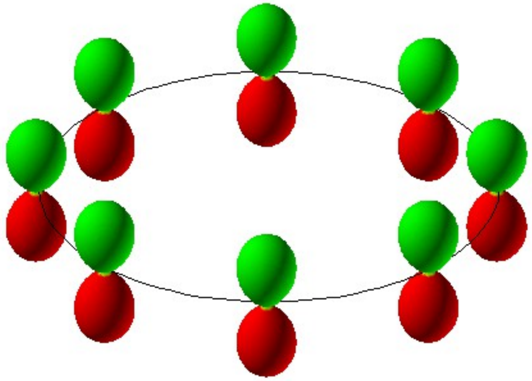


Topological excitation in polar state



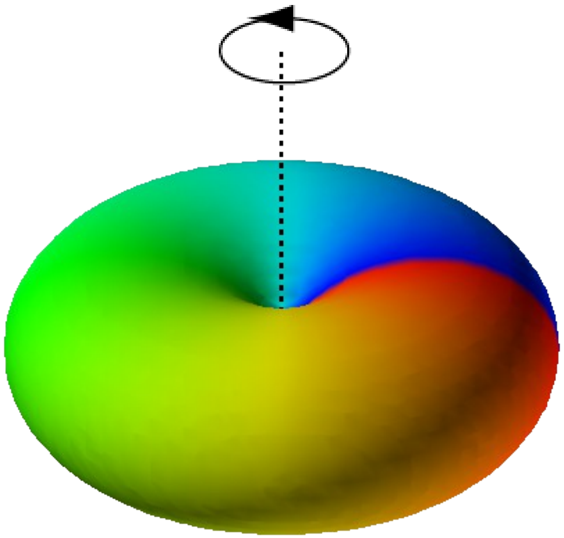
vortex

$$\frac{G}{H} \simeq \frac{U(1)_\varphi \times S_F^2}{(\mathbb{Z}_2)_{\varphi+F}}$$



not vortex

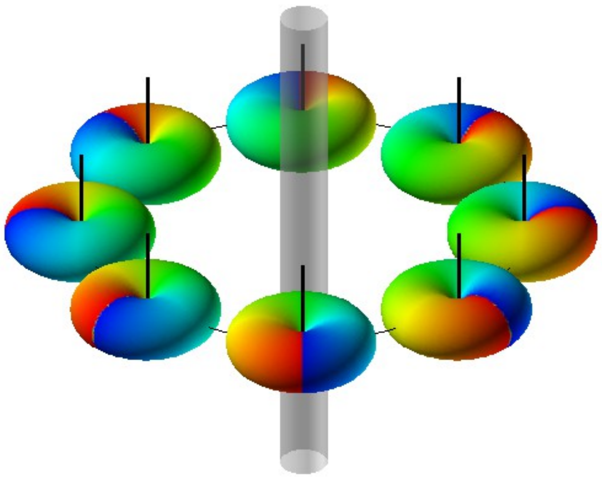
Topological excitation in Ferromagnetic state



$$\frac{G}{H} \simeq \frac{U(1)_\varphi \times SO(3)_F}{(U(1))_{\varphi+F}}$$

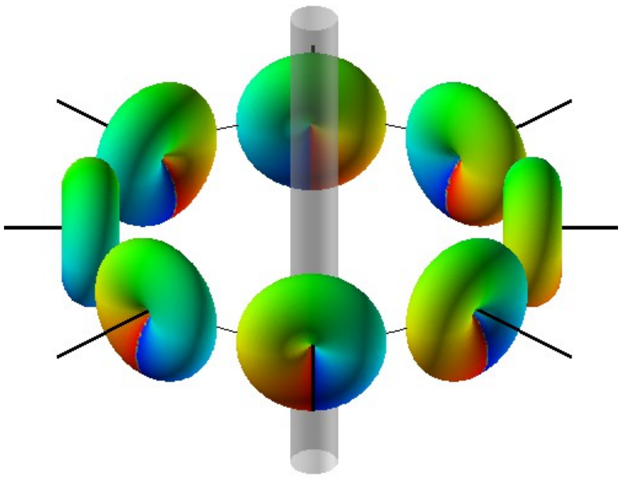
$$\simeq SO(3)_{\varphi+F}$$

Gauge vortex



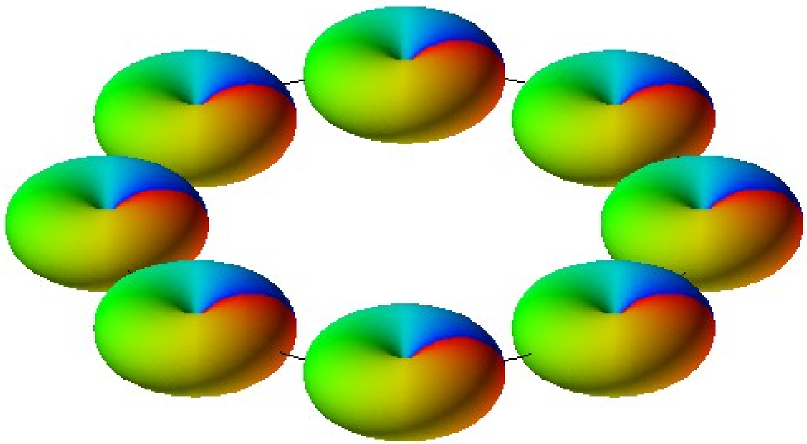
Continuous transformation

Spin vortex



Topological excitation in Ferromagnetic state

Doubly winding state is no longer topological excitation

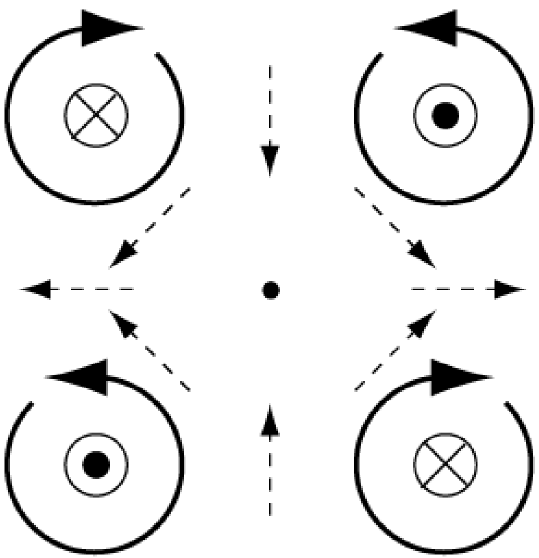


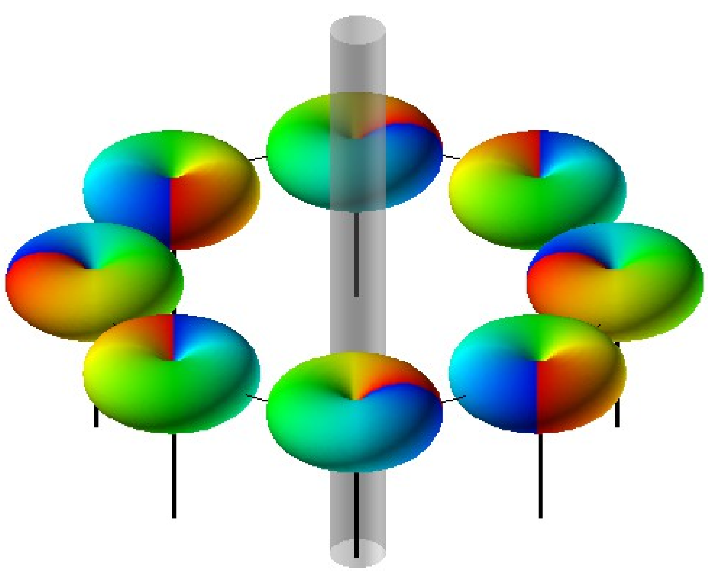
$$\pi_1[SO(3)] \simeq \mathbb{Z}_2$$

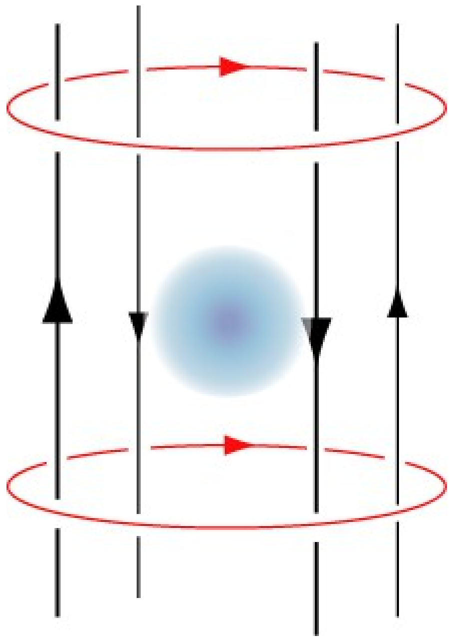
T. Ishoshima, et al. PRA **61**, 063610 (1999)

Topological excitation in Ferromagnetic state

Creation of doubly winding state from zero winding



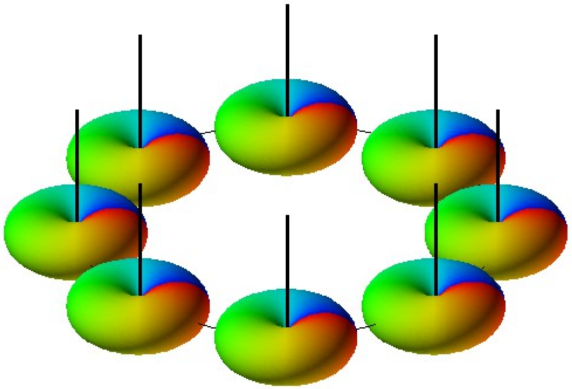




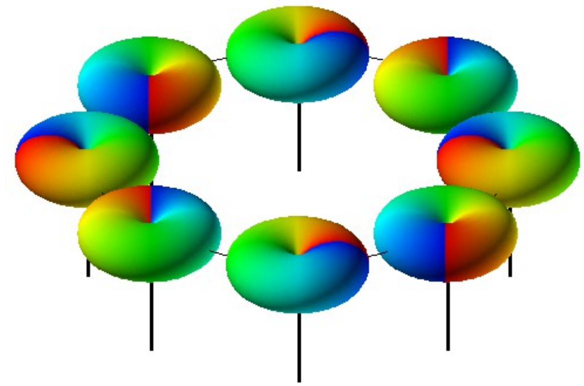
Adiabatic change of quadratic magnetic field

Topological excitation in Ferromagnetic state

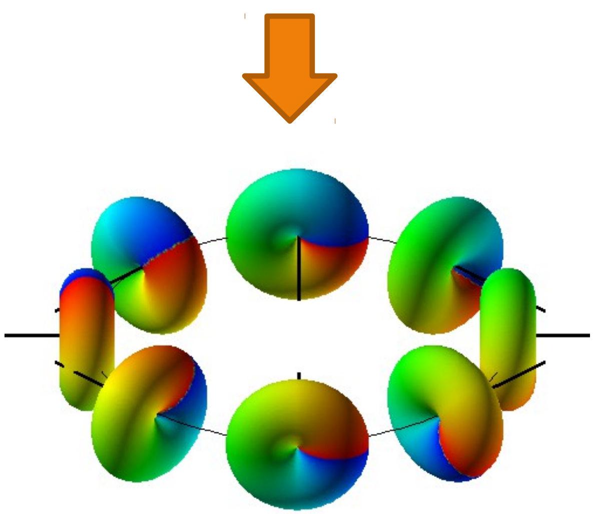
Creation of doubly winding state from zero winding

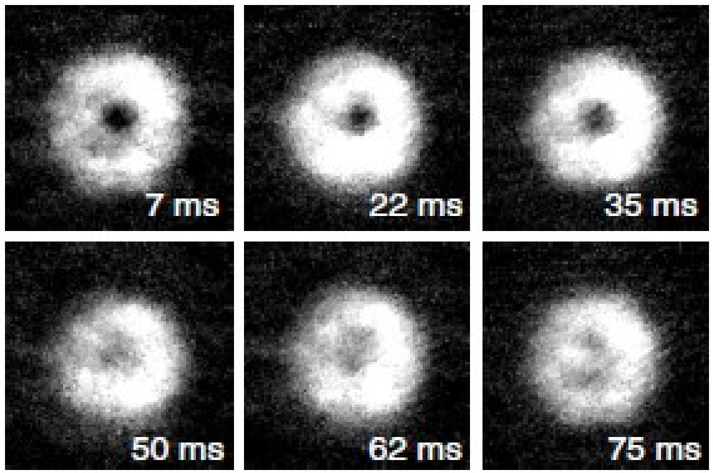


$$\begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}$$



$$\begin{pmatrix} 0 \\ 0 \\ e^{-2i\theta} \end{pmatrix}$$





Y. Shin, et al. PRL **93**, 160406 (2004)

Spin-2 case

$$H = \int dx \left[-\Psi_m^* \frac{\hbar^2}{2M} \nabla^2 \Psi_m + \frac{c_0}{2}\rho^2 + \frac{c_1}{2}\boldsymbol{F}^2 + \frac{c_2}{2}|A_{20}|^2 \right]$$

$$A_{00}(\boldsymbol{x}) = 2\Psi_2(\boldsymbol{x})\Psi_{-2}(\boldsymbol{x}) - 2\Psi_1(\boldsymbol{x})\Psi_{-1}(\boldsymbol{x}) + \Psi_0(\boldsymbol{x})^2$$

Singlet-pair amplitude

Spin-2 case

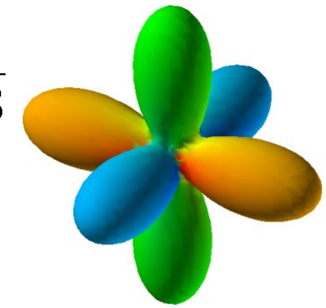
$$H=\int d\boldsymbol{x}\left[-\Psi_m^*\frac{\hbar^2}{2M}\nabla^2\Psi_m+\frac{c_0}{2}\rho^2+\frac{c_1}{2}\boldsymbol{F}^2+\frac{c_2}{2}|A_{20}|^2\right]$$

c_1

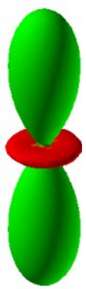
Cyclic:

$\Psi_\mathrm{C}=(1,0,0,\sqrt{2},1)^T/\sqrt{3}$

$\frac{U(1)_\varphi\times SO(3)_F}{(T)_{\varphi+F}}$



Uniaxial Nematic:



$\Psi_\mathrm{U}=(0,0,1,0,0)^T$

$U(1)_\varphi\times\frac{S_F^2}{(\mathbb{Z}_2)_F}$

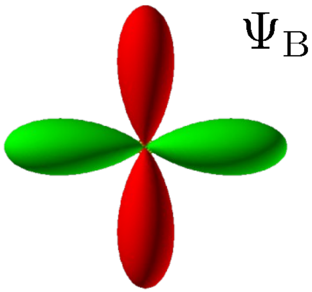
^{87}Rb

c_2

Biaxial Nematic:

$\Psi_\mathrm{B}=(1,0,0,0,1)^T/\sqrt{2}$

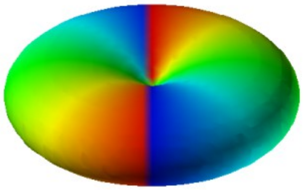
$\frac{U(1)_\varphi\times SO(3)_F}{(D_4)_{\varphi+F}}$



Ferromagnetic:

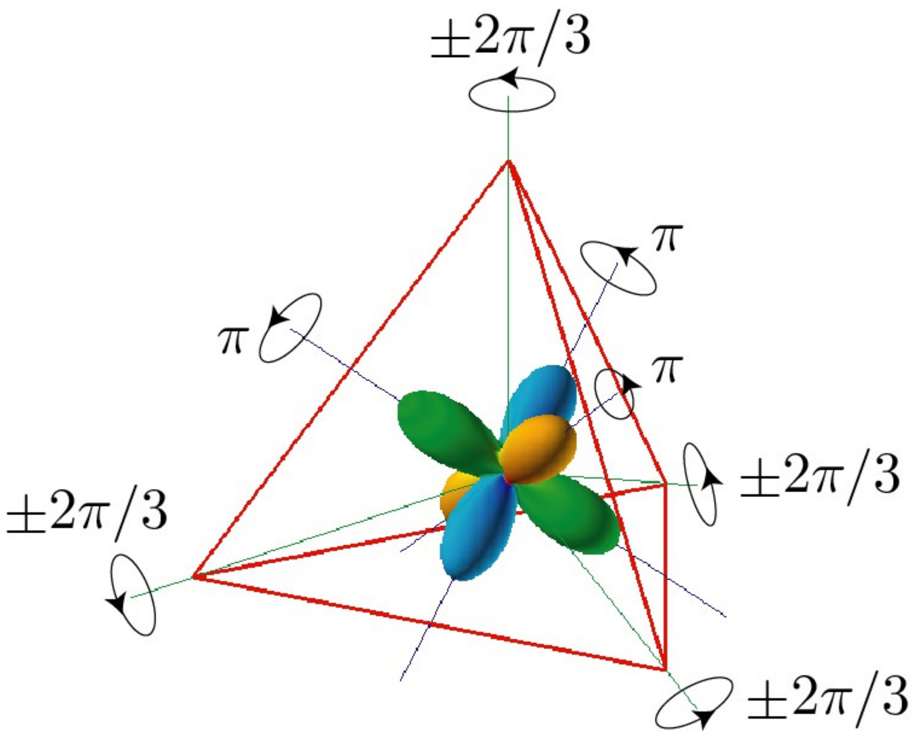
$\Psi_\mathrm{F}=(1,0,0,0,0)^T$

$\frac{SO(3)_{\varphi+F}}{(\mathbb{Z}_2)_{\varphi+F}}$



$c_2=4c_1$

Topological excitation in cyclic state



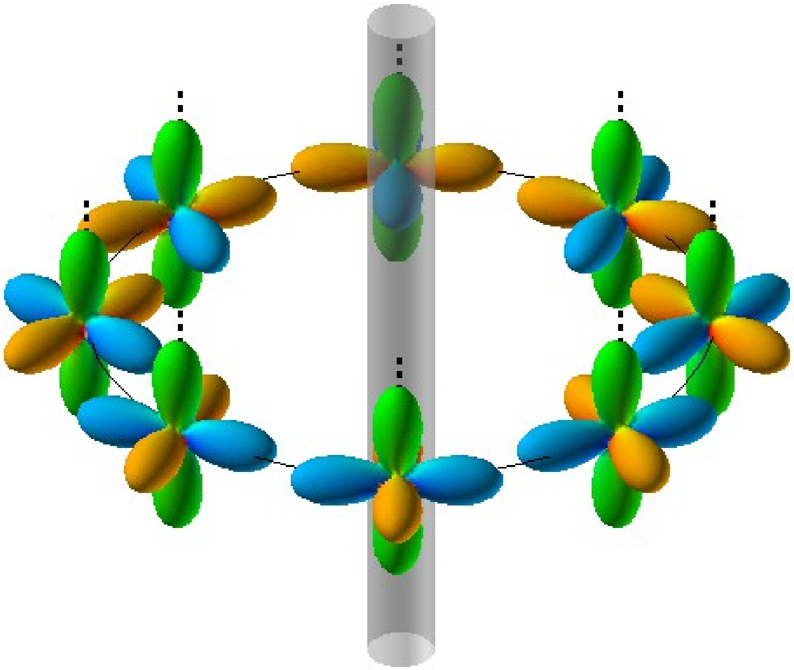
$$\frac{U(1)_\varphi \times SO(3)_F}{(T)_{\varphi+F}}$$

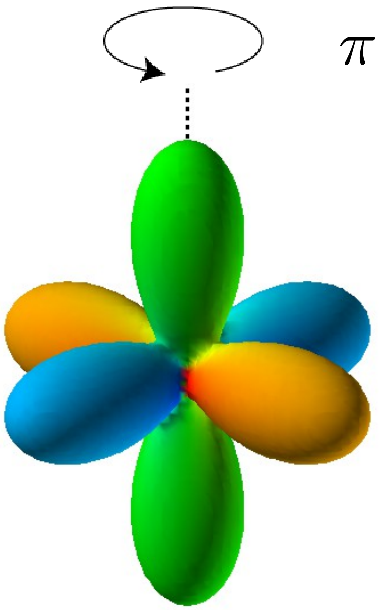
Topological excitations can be labeled by12 rotations keeping tetrahedron invariant

Non-Abelian topological excitation!

Topological excitation in cyclic state

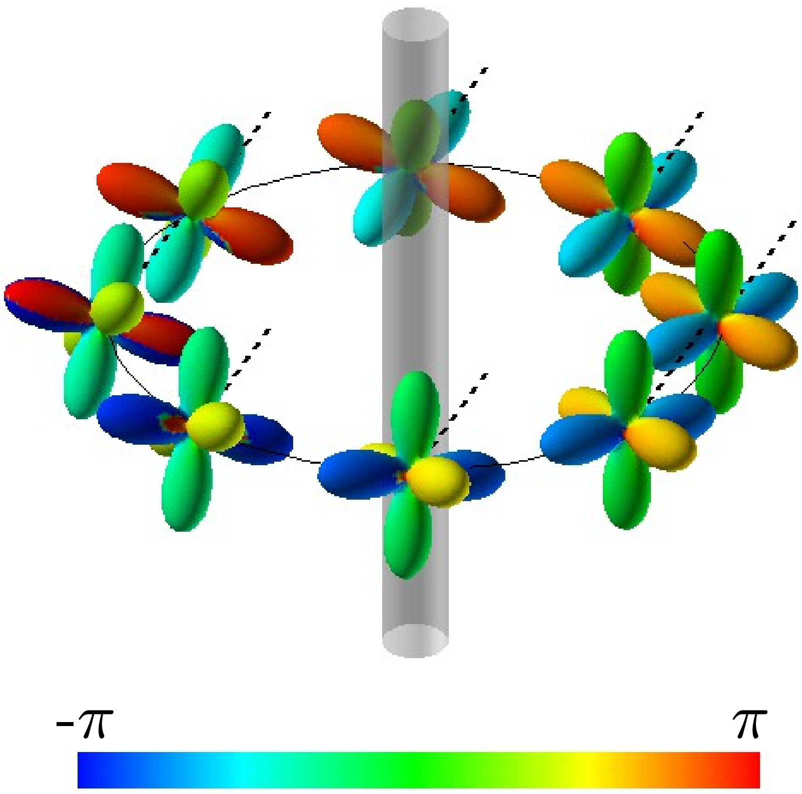
1/2-spin vortex

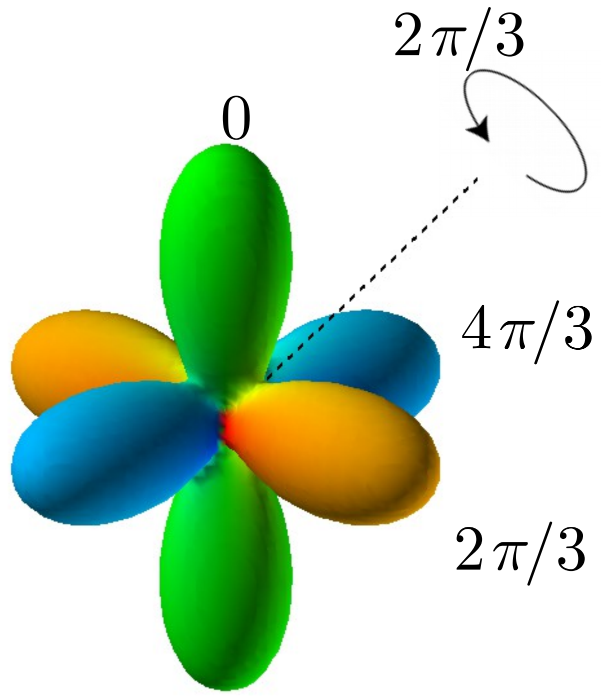




Topological excitation in cyclic state

1/3 vortex





Dynamics of topological excitations

Collision of topological excitations

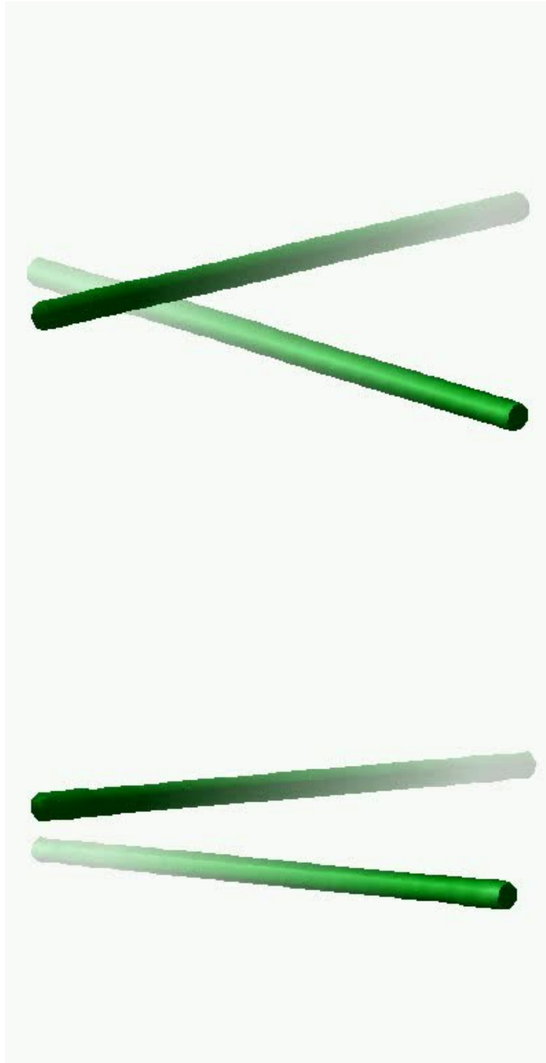
Quantum turbulence

Collision of vortices with non-commutative charge forms a new “**rung**” vortex connecting two vortices

Reconnection

Abelian excitation

Passing



Non-Abelian excitation

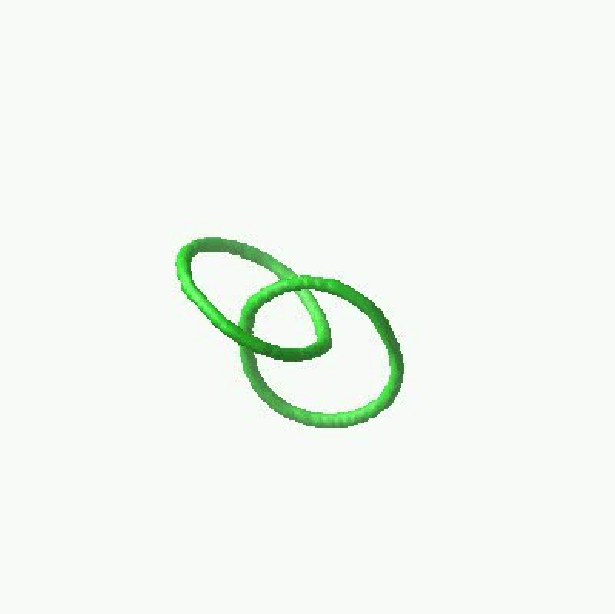
Rung vortex

M. Kobayashi, et al. PRL **103**, 115301(2009)

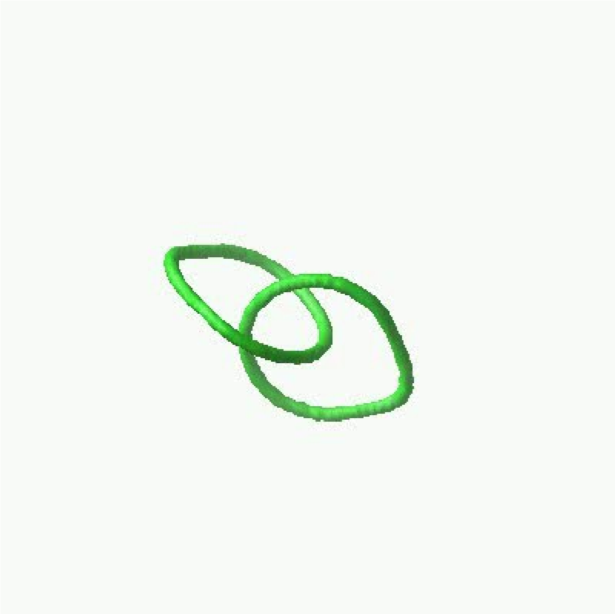
Collision dynamics of topological excitations

Abelian excitation

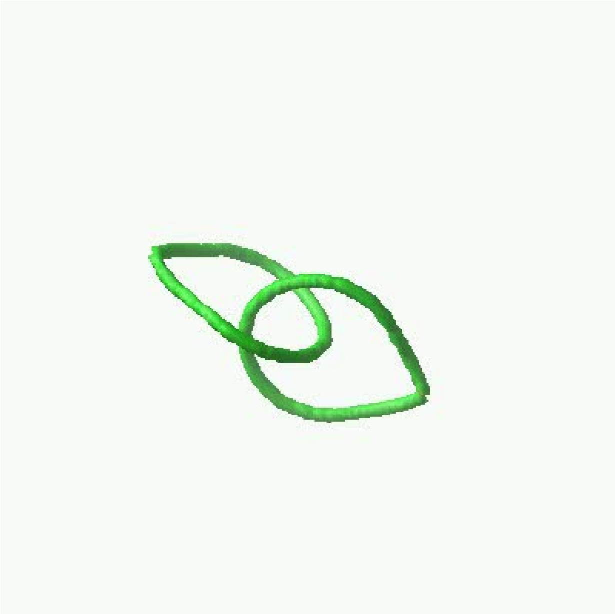
Non-Abelian excitation



Large ring



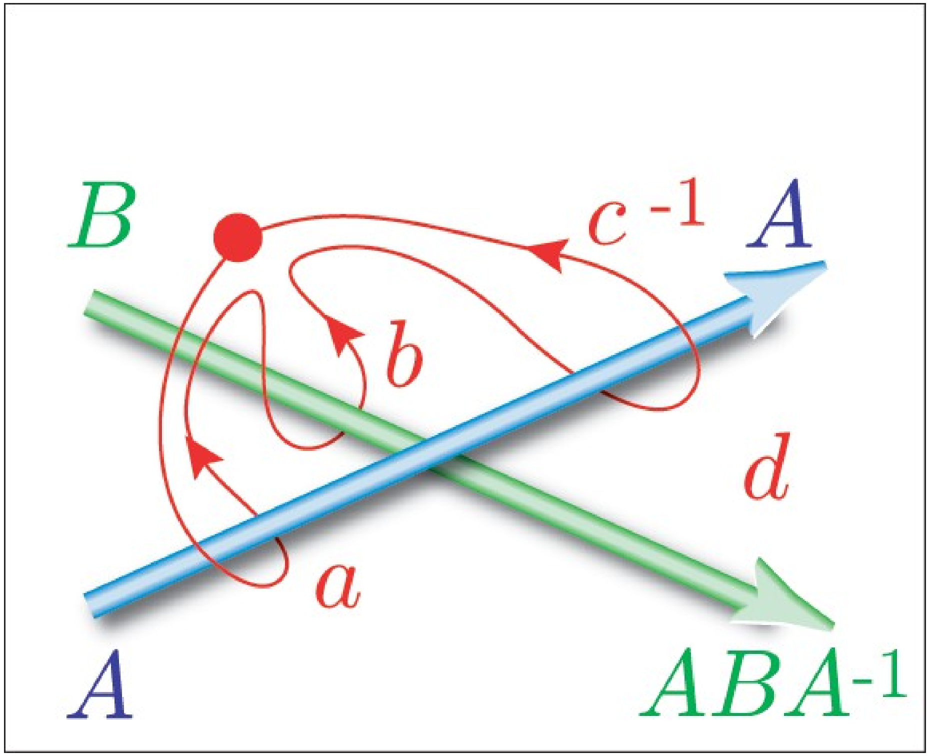
Unraveling of link



New excitation

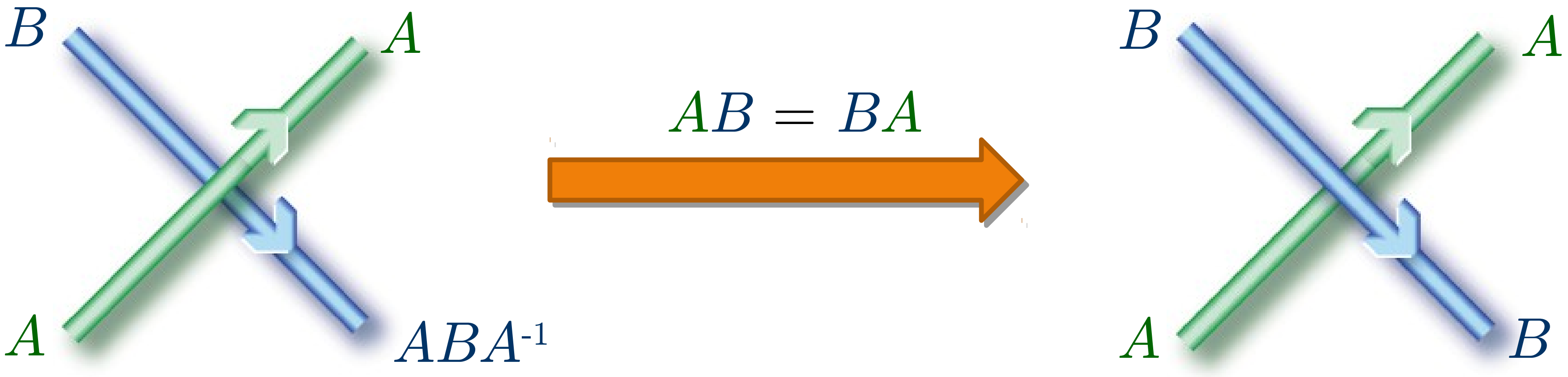
Linked non-Abelian excitations cannot unravel because of the formation of the new excitation.

Topological charge of topological excitation



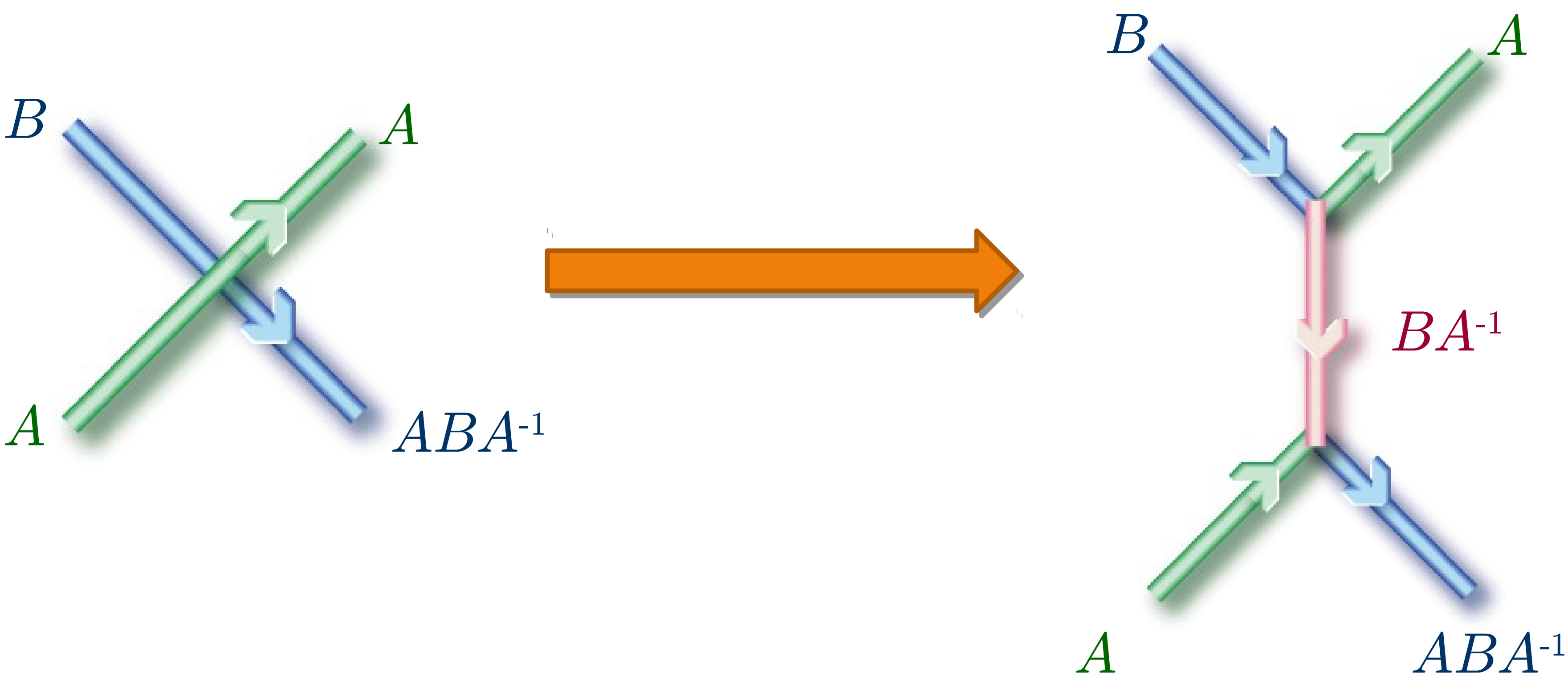
Topological invariant of excitations can be fixed by a closed path encircling the excitations

Collision of Vortex



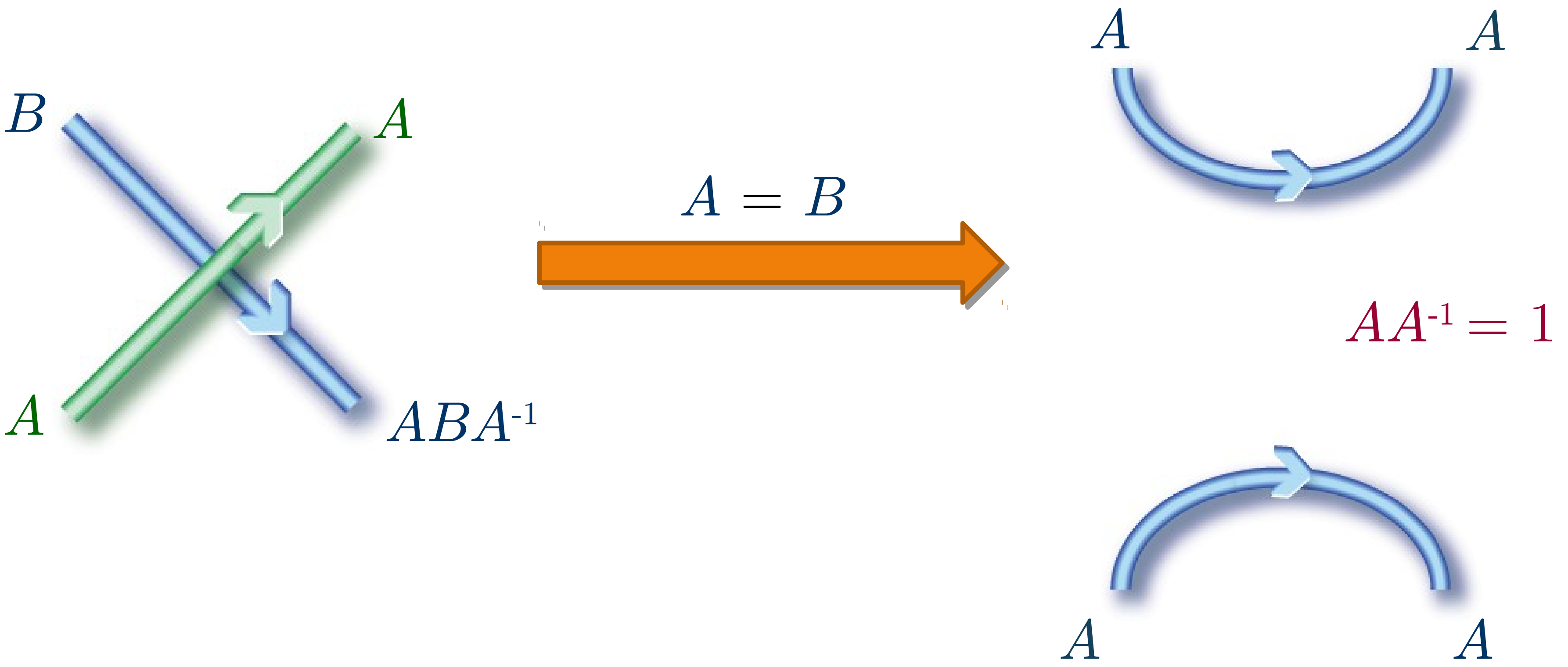
Passing dynamics is possible for Abelian case

Collision of Vortex



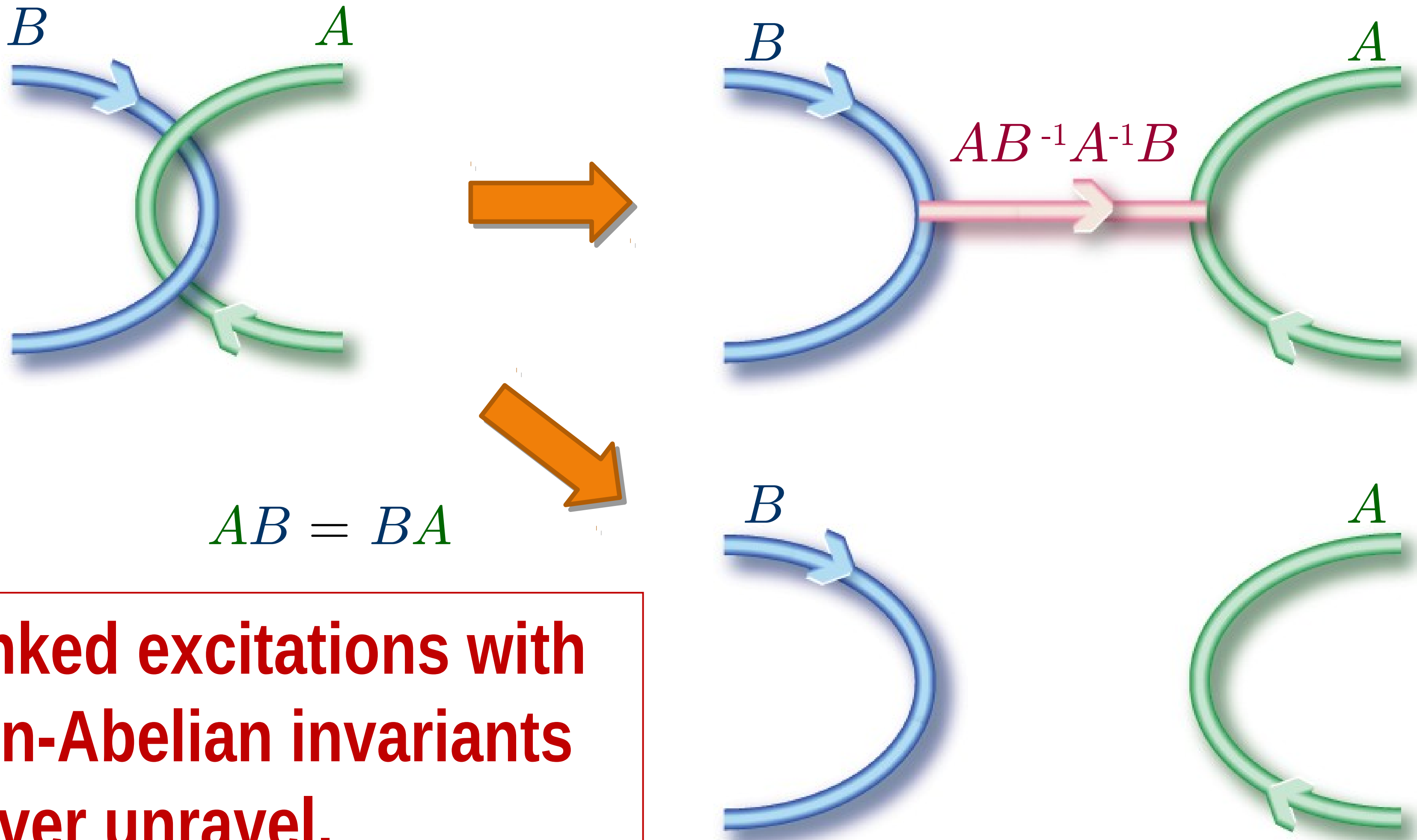
Rung BA^{-1} is formed through the collision.

Collision of Vortex



Rung disappears for the same charge resulting reconnection.

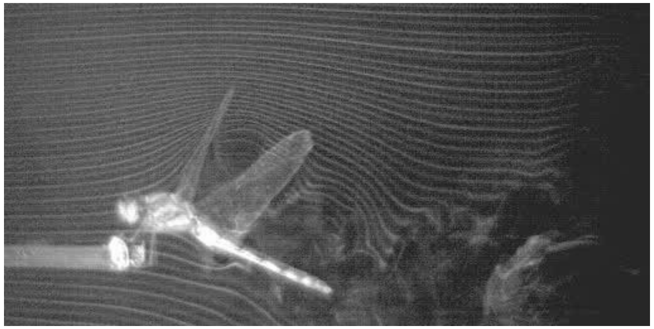
Linked Vortex Rings



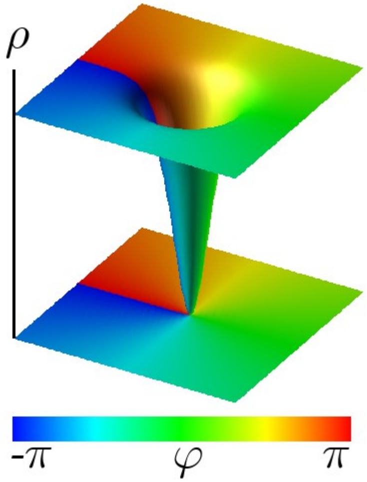
Linked excitations with non-Abelian invariants never unravel.

Quantum turbulence

Turbulent state of BEC (quantum turbulence) can be a model to understand the relation between turbulence and vortices



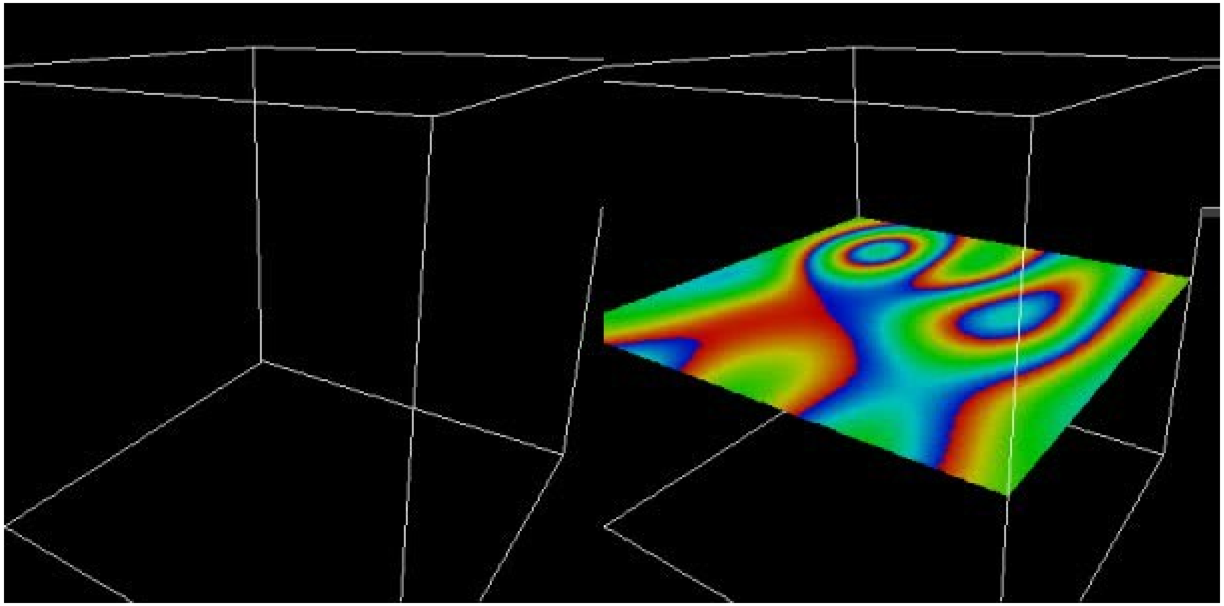
Vortices in viscous fluid cannot be defined clearly



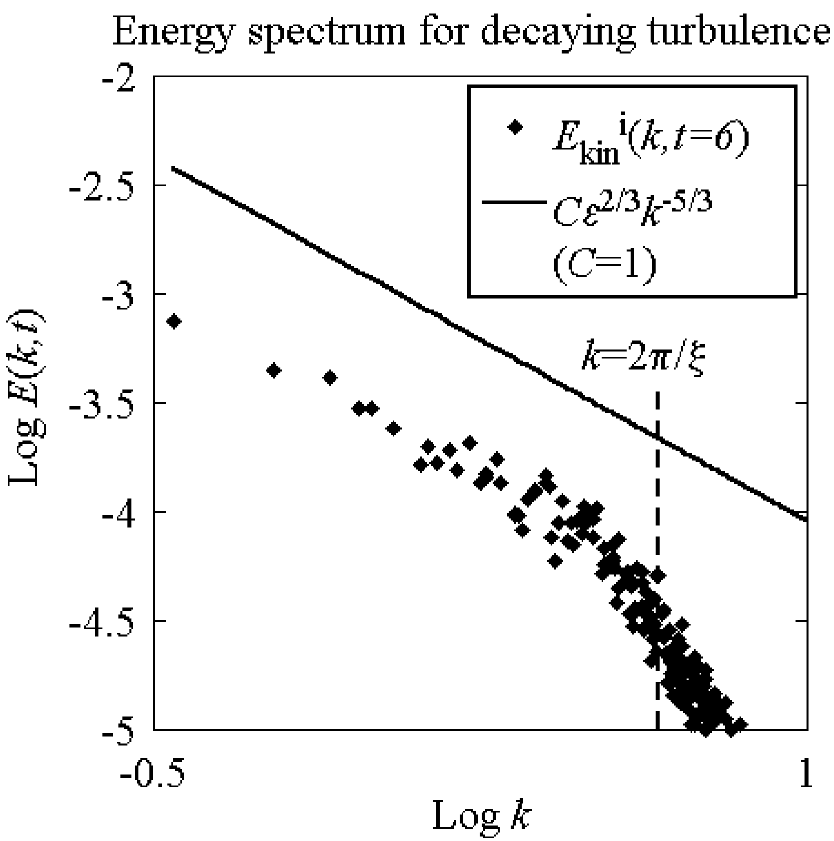
Vortices in BEC can be defined as quantized vortices which are topological excitations

Generation of turbulence

Vortices 2D plane of the phase



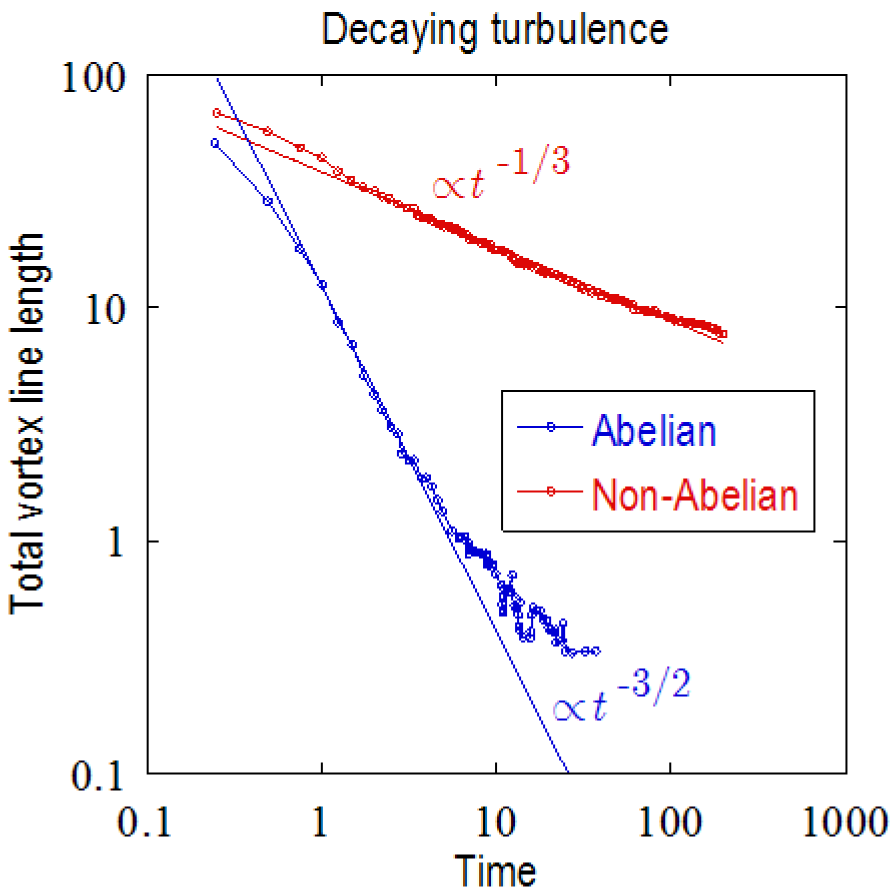
Energy spectrum of quantum turbulence

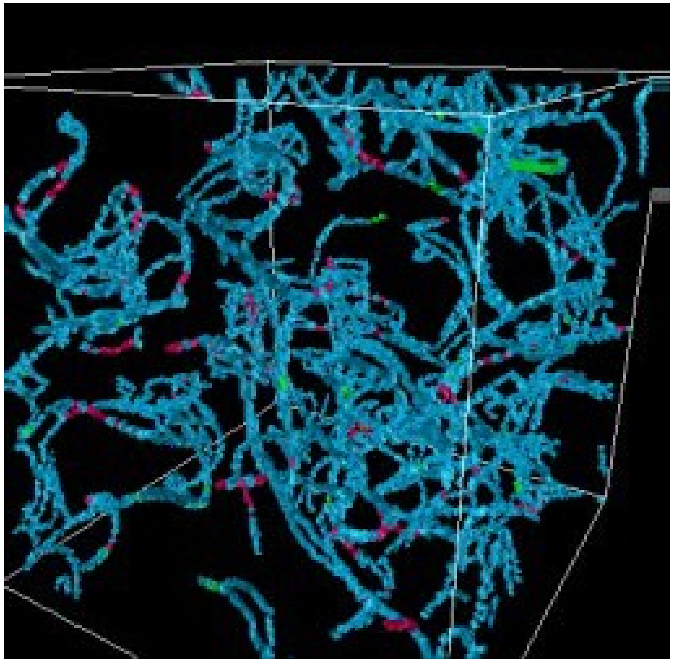


Quantum turbulence shows the Kolmogorov law which is one of the most important statistical law of turbulence

M. Kobayashi, et al. PRL **94**, 065302 (2005)

Non-Abelian quantum turbulence





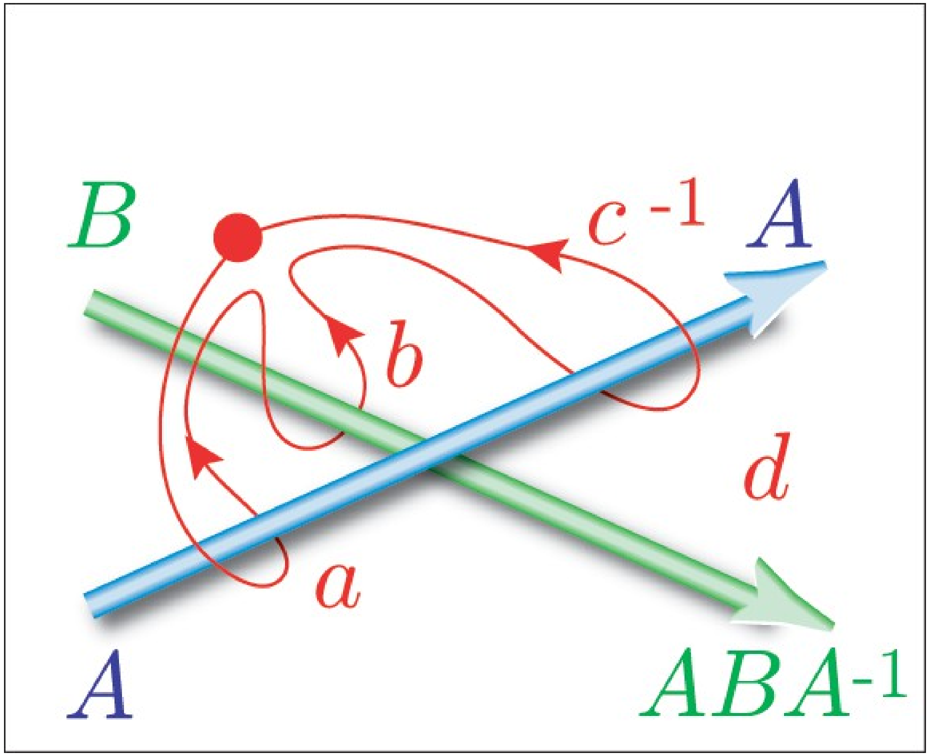
Turbulent behavior is strongly affected by the topology

M. Kobayashi, et al. in press

Summary

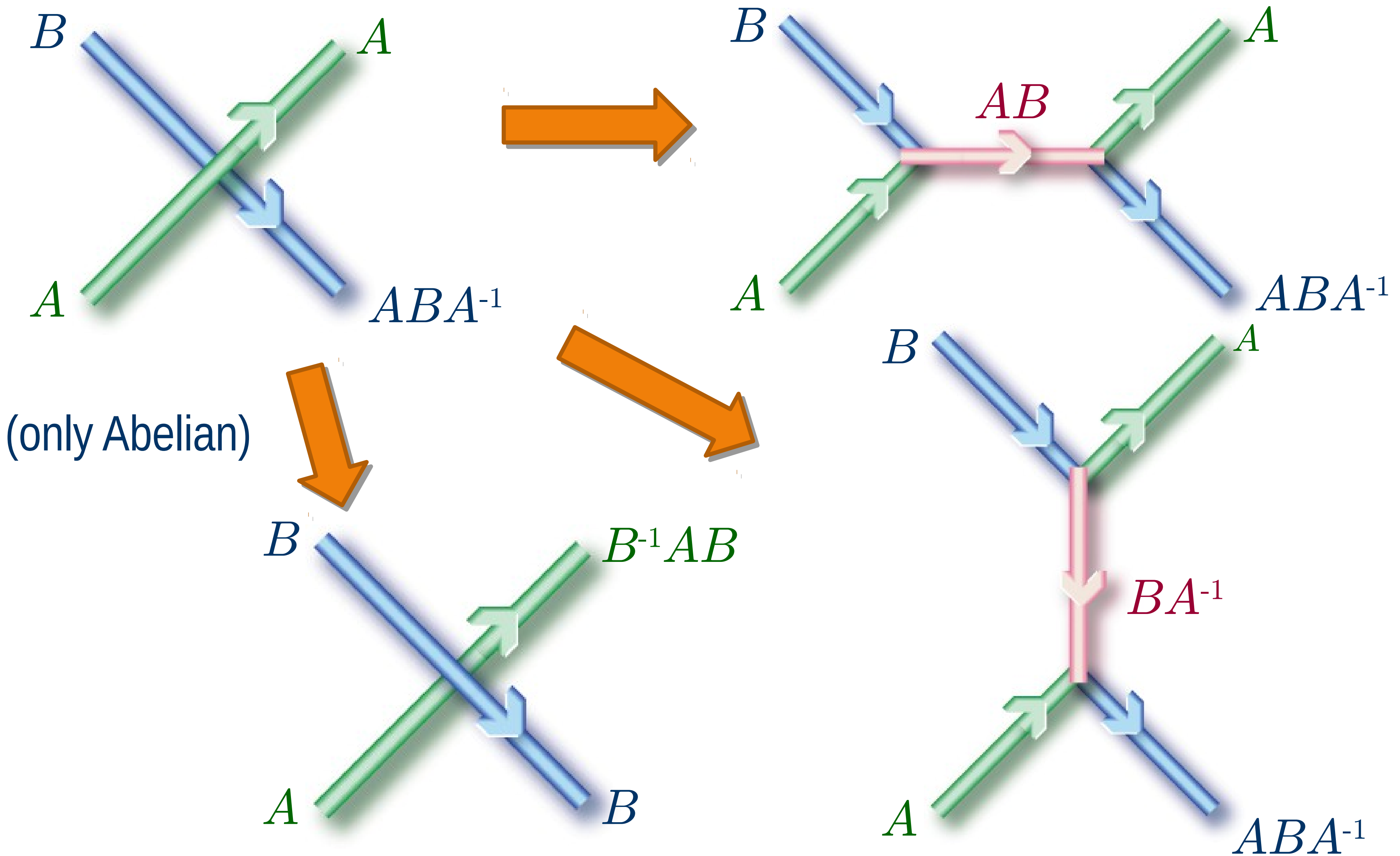
- In BEC, various kinds of topological excitations can be realized.
- Dynamics of topological excitations are affected by the order-parameter manifold and can dominate the nature of the system.

Topological charge of topological excitation

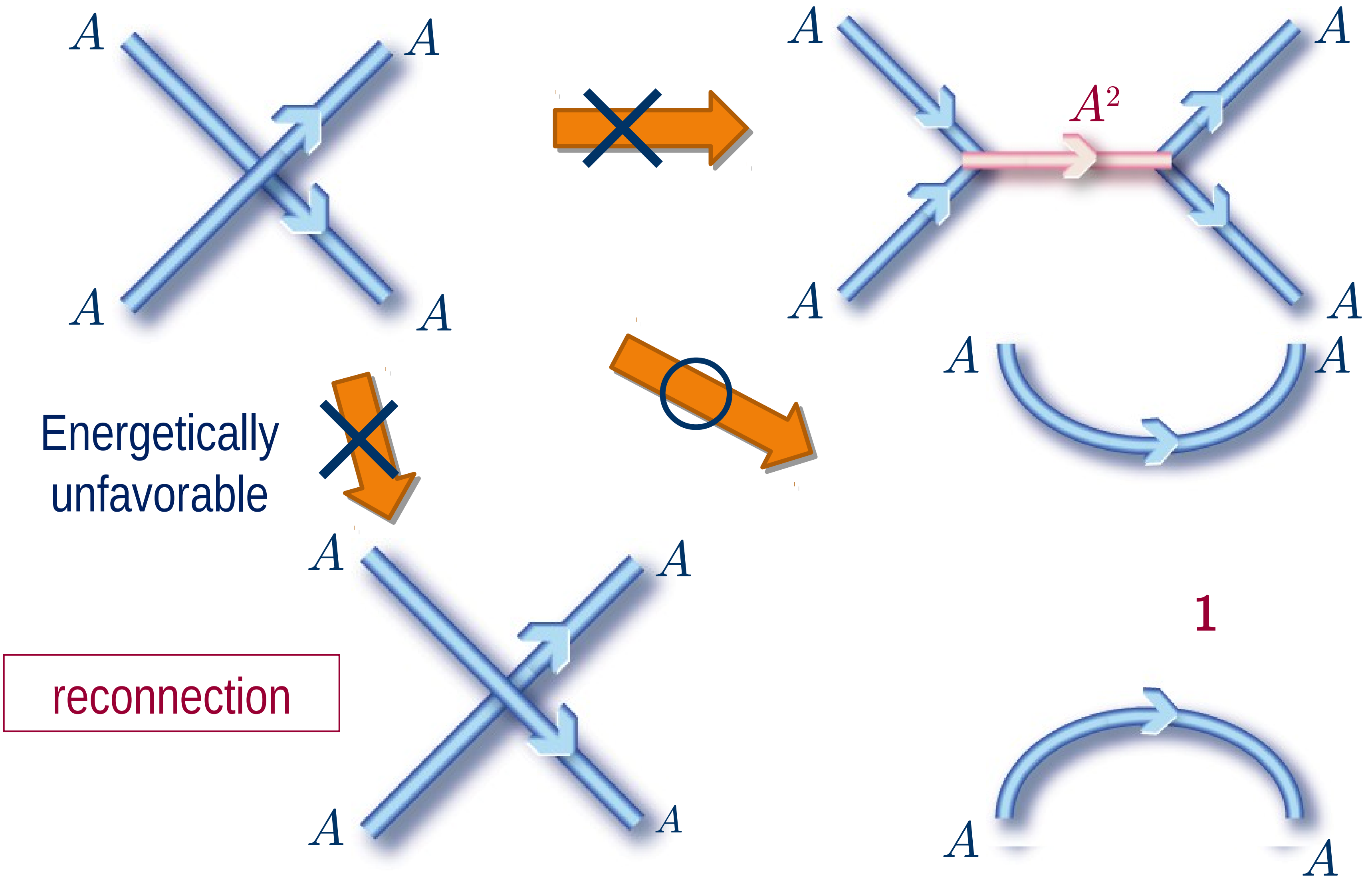


Topological charge of vortex can be fixed by a closed path encircling the vortex

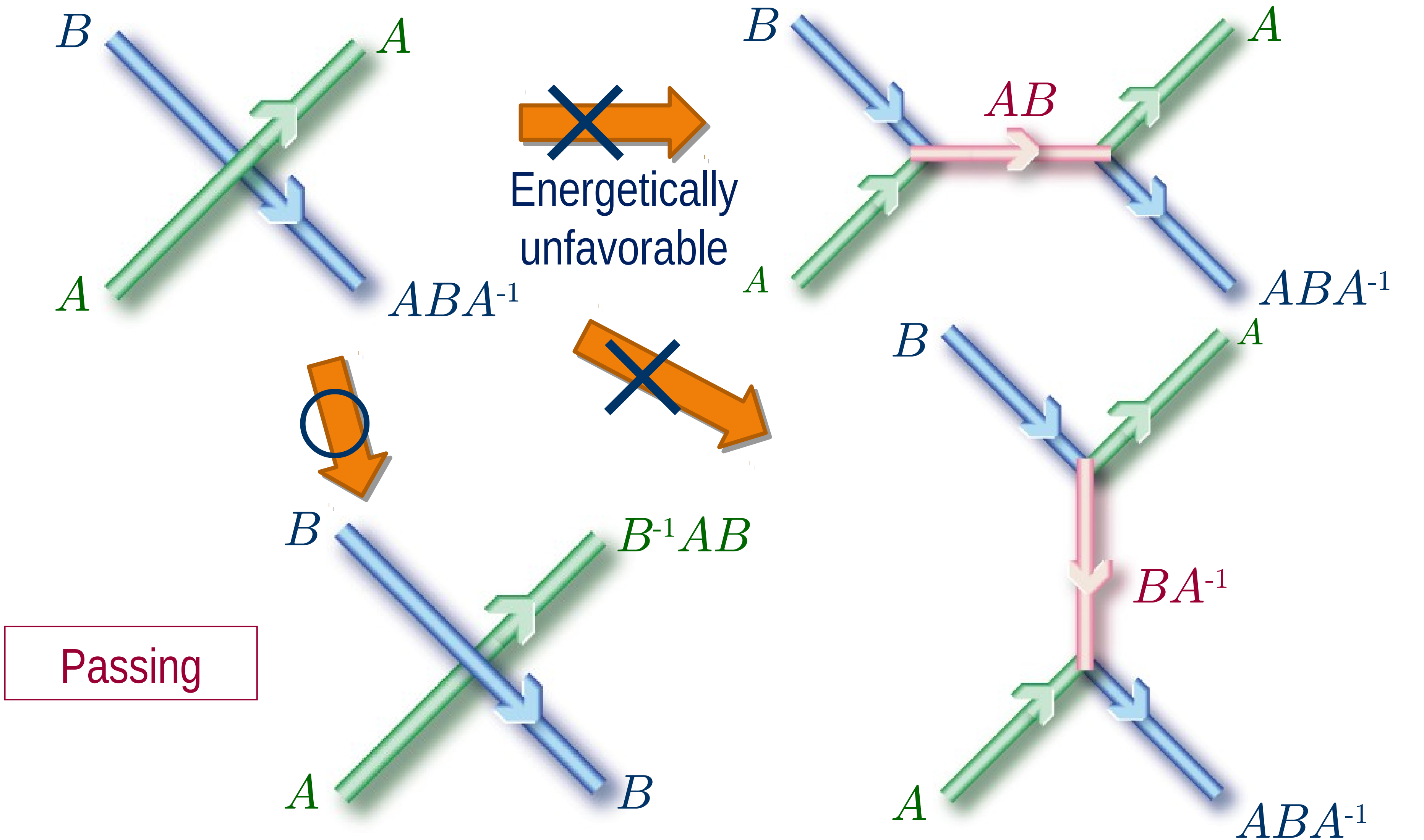
Collision of Vortices



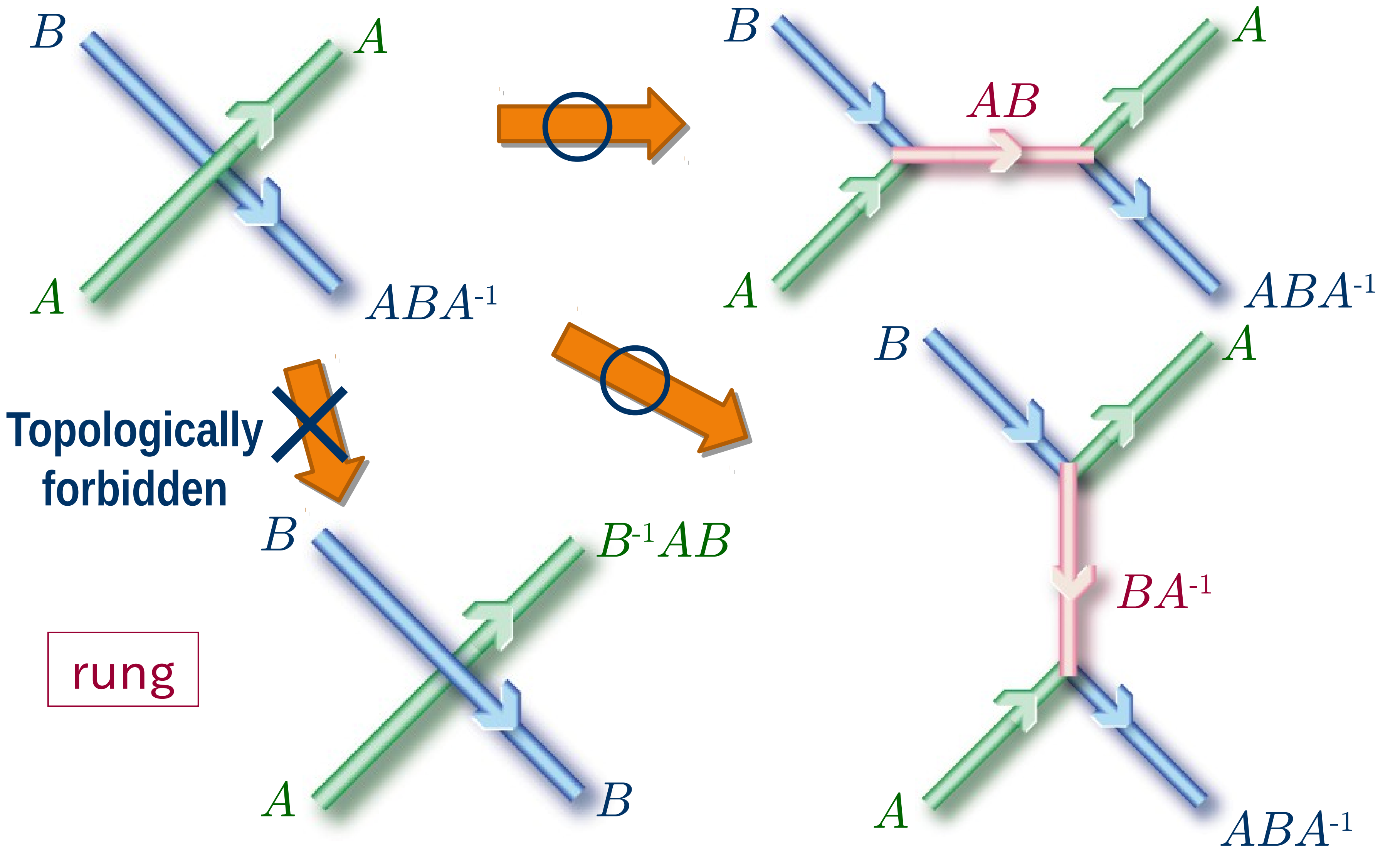
Collision of Same Vortices



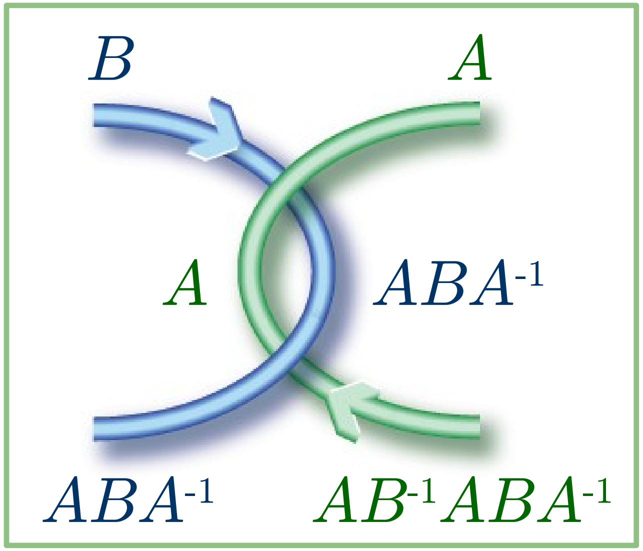
Collision of Different Commutative Vortices

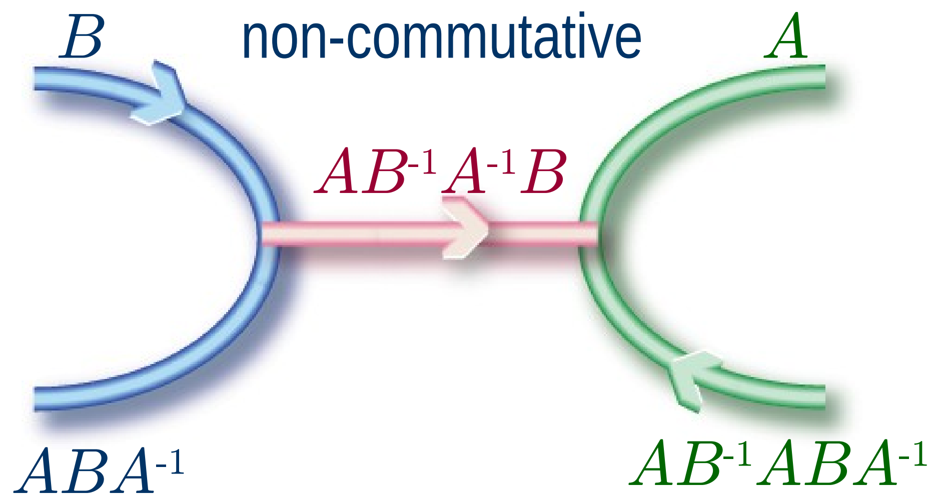


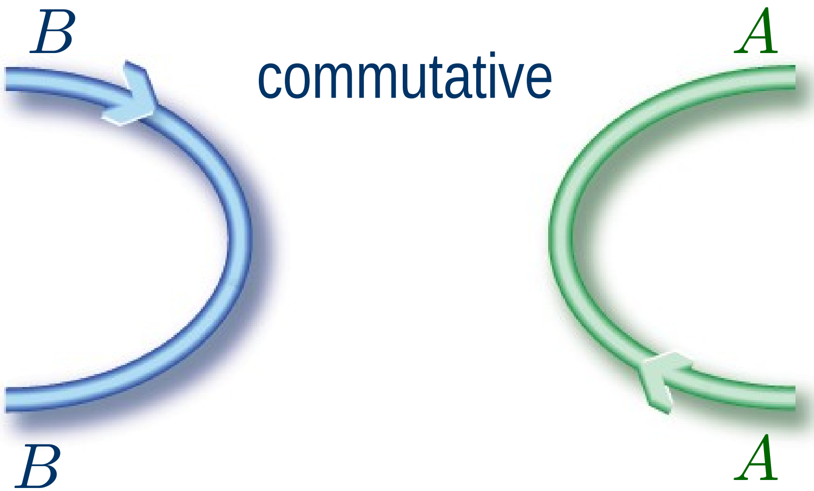
Collision of Different Non-commutative Vortices



Linked Vortices



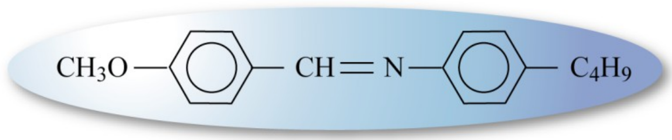


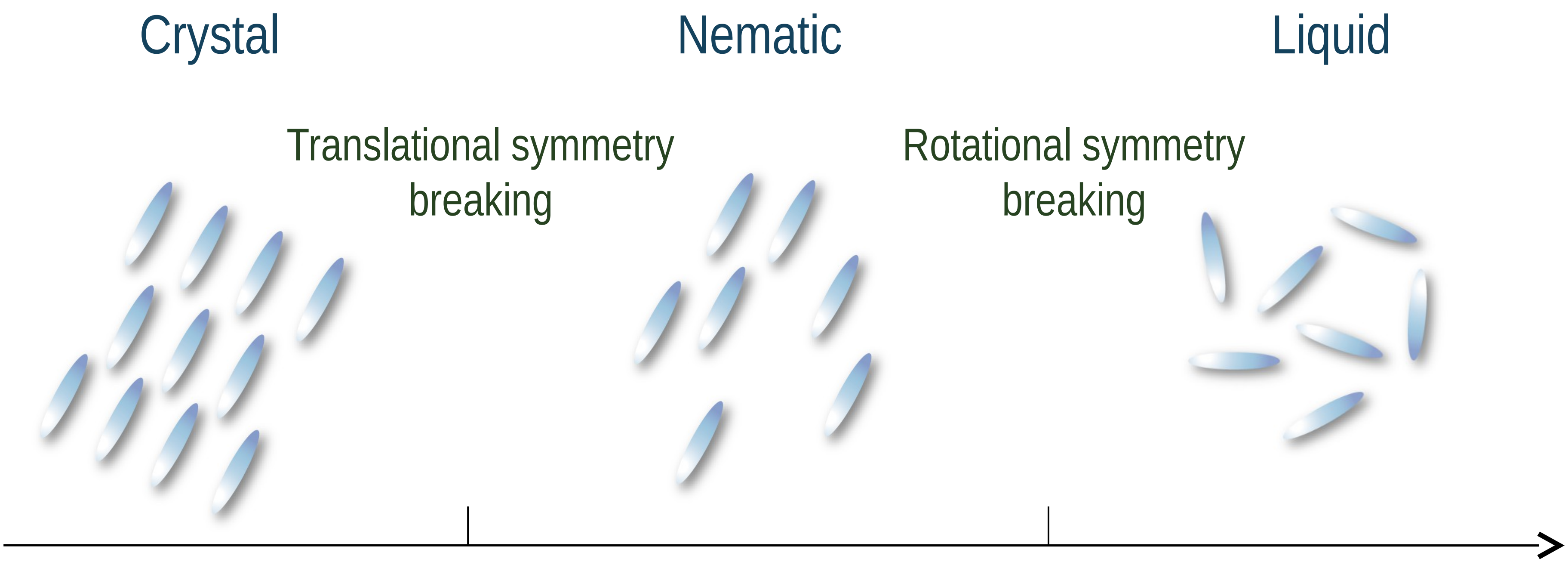


Linked vortices cannot untangle

What is topological excitations?

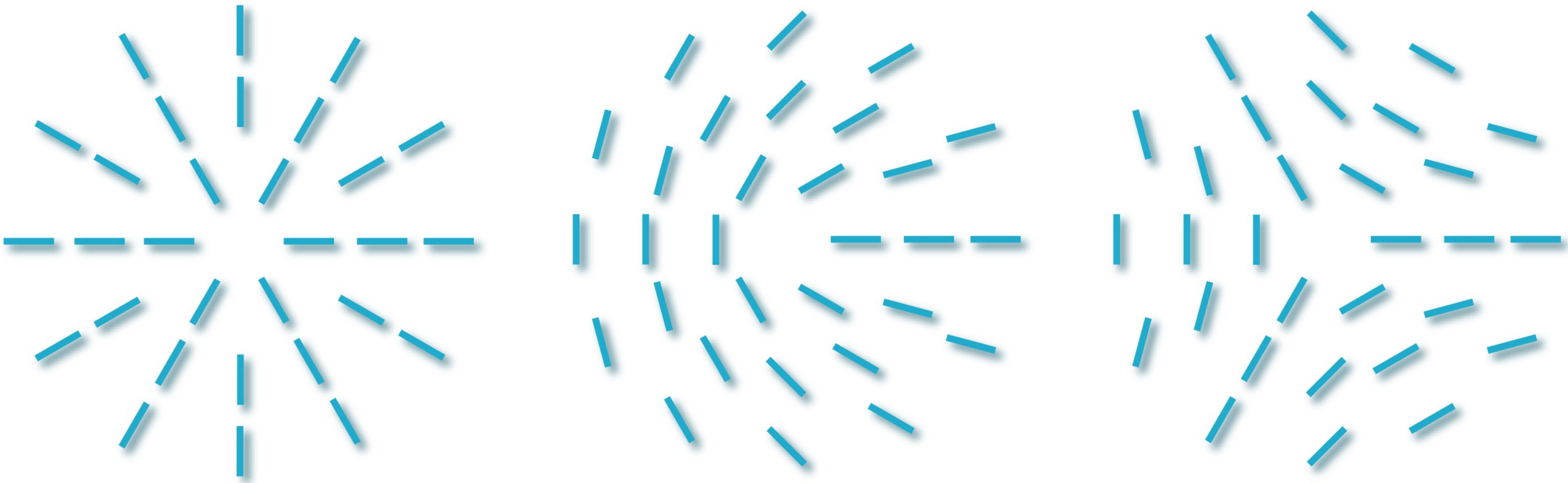
Nematic liquid crystal





What is topological excitations?

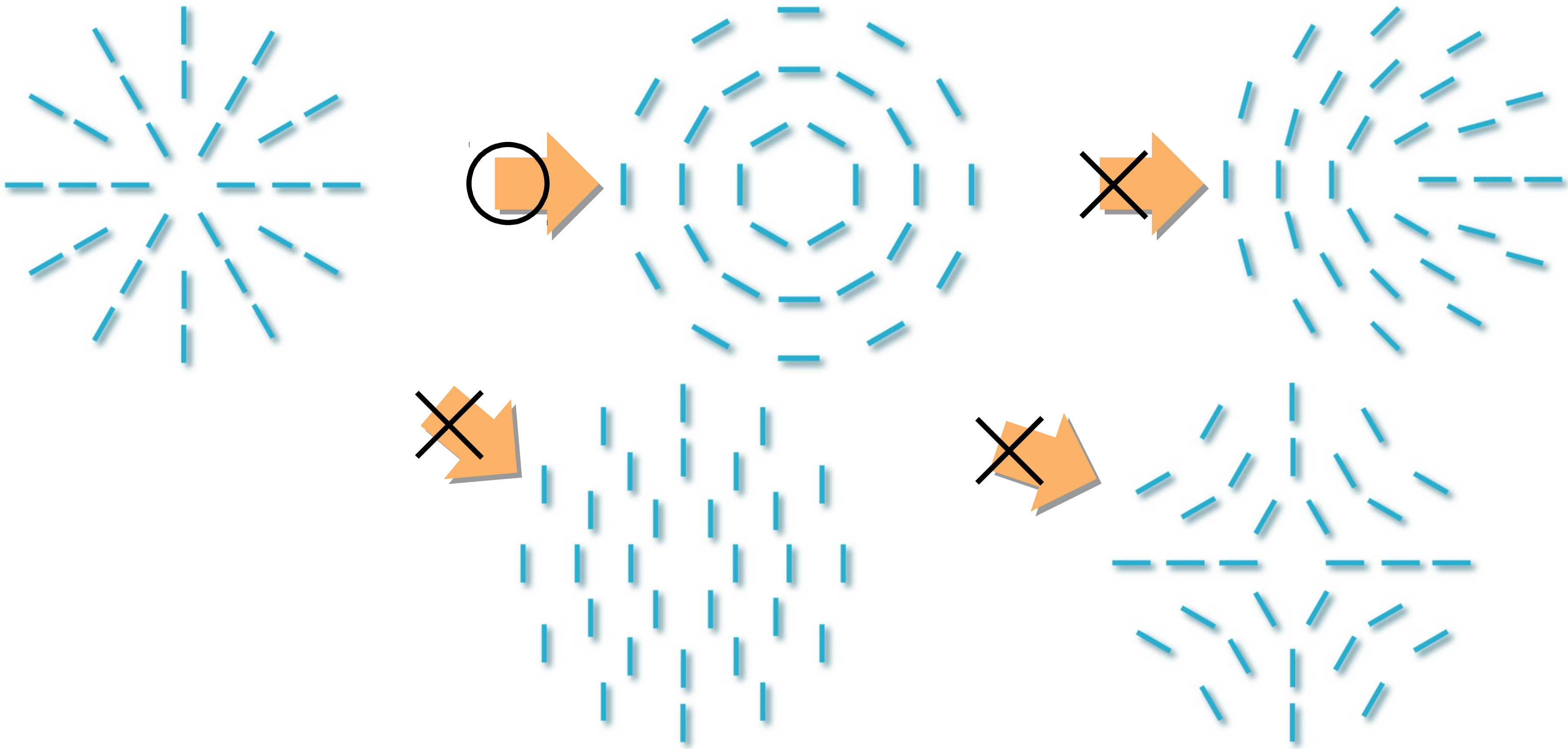
Topological excitations in nematic liquid crystal



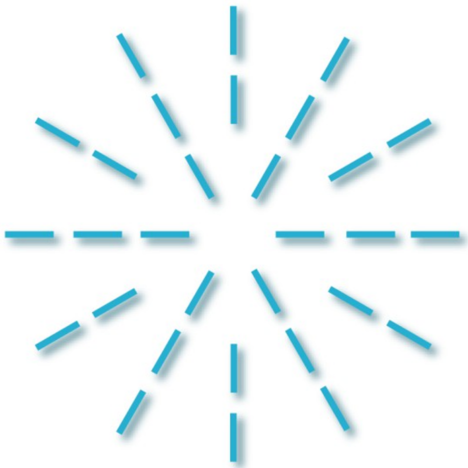
Topological excitation related to rotational symmetry breaking

What is topological excitations?

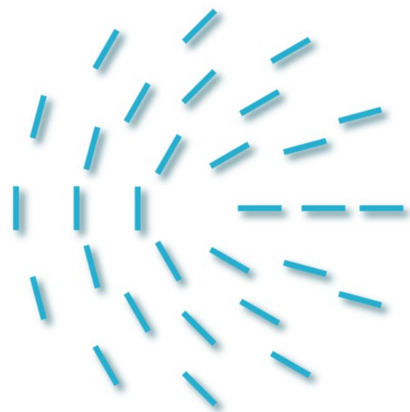
States with topological excitations cannot be continuously transformed to states without topological excitations



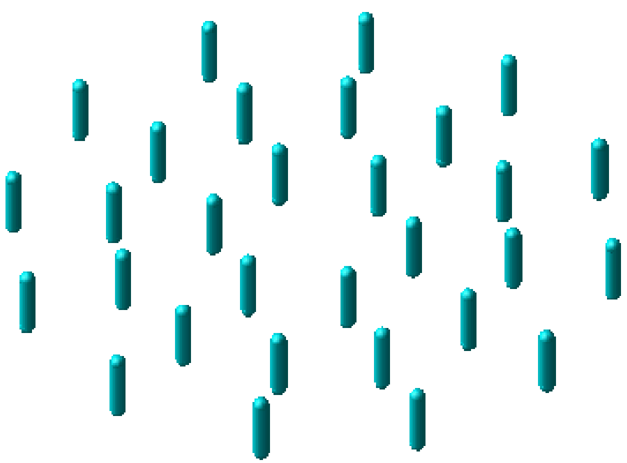
What is topological excitations?



This is topological excitation in 2D system but not topological excitation in 3D system

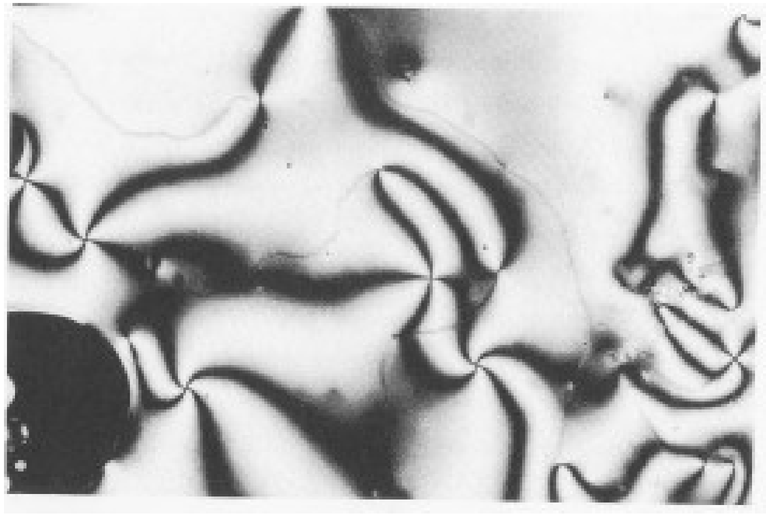


This is always topological excitation

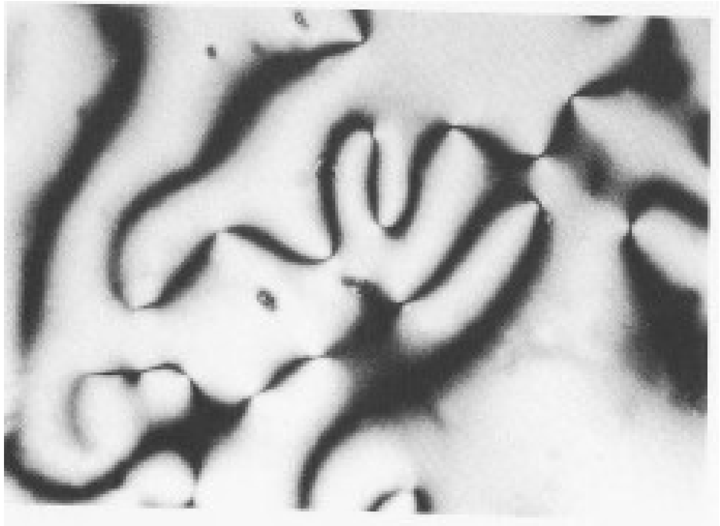


Characteristics of topological excitations strongly depend on the internal degrees of freedom (topology) of the system

Observation of topological excitations in nematic liquid crystal

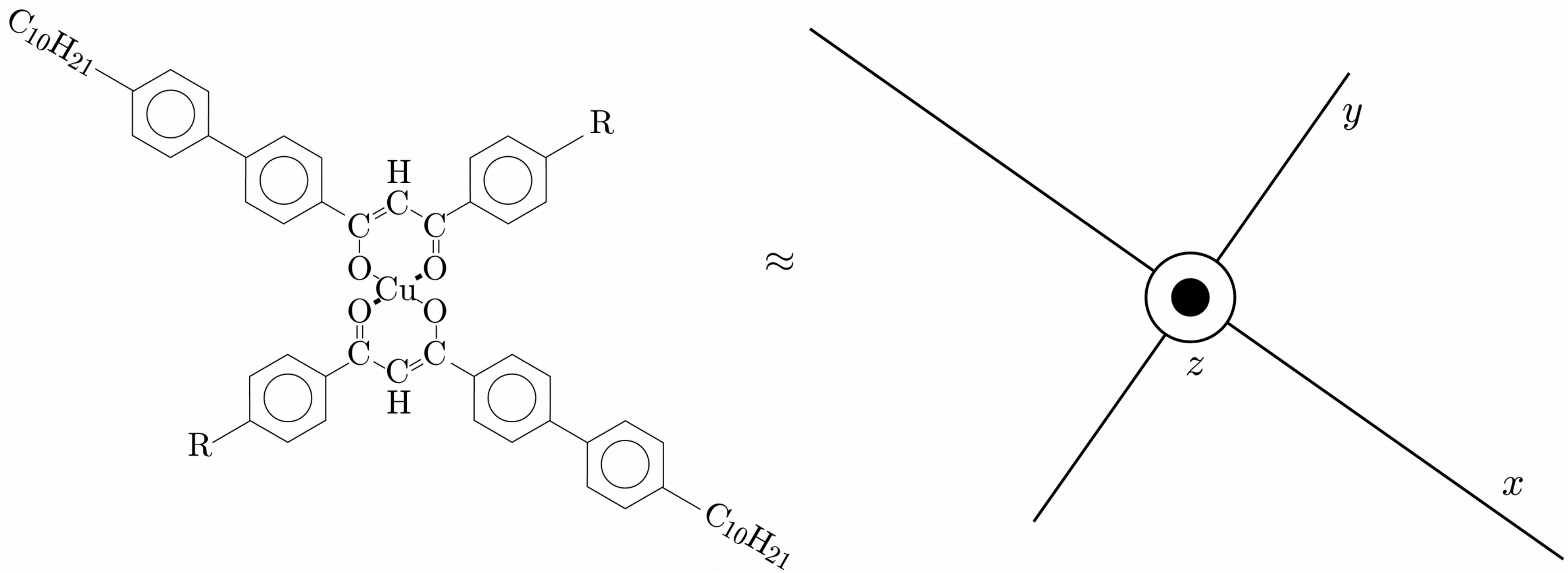


near surface



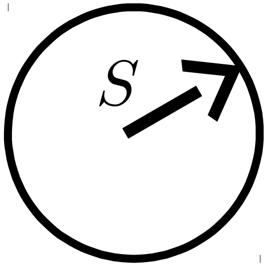
far from surface

Biaxial nematic liquid crystal



Topological excitations and homotopy

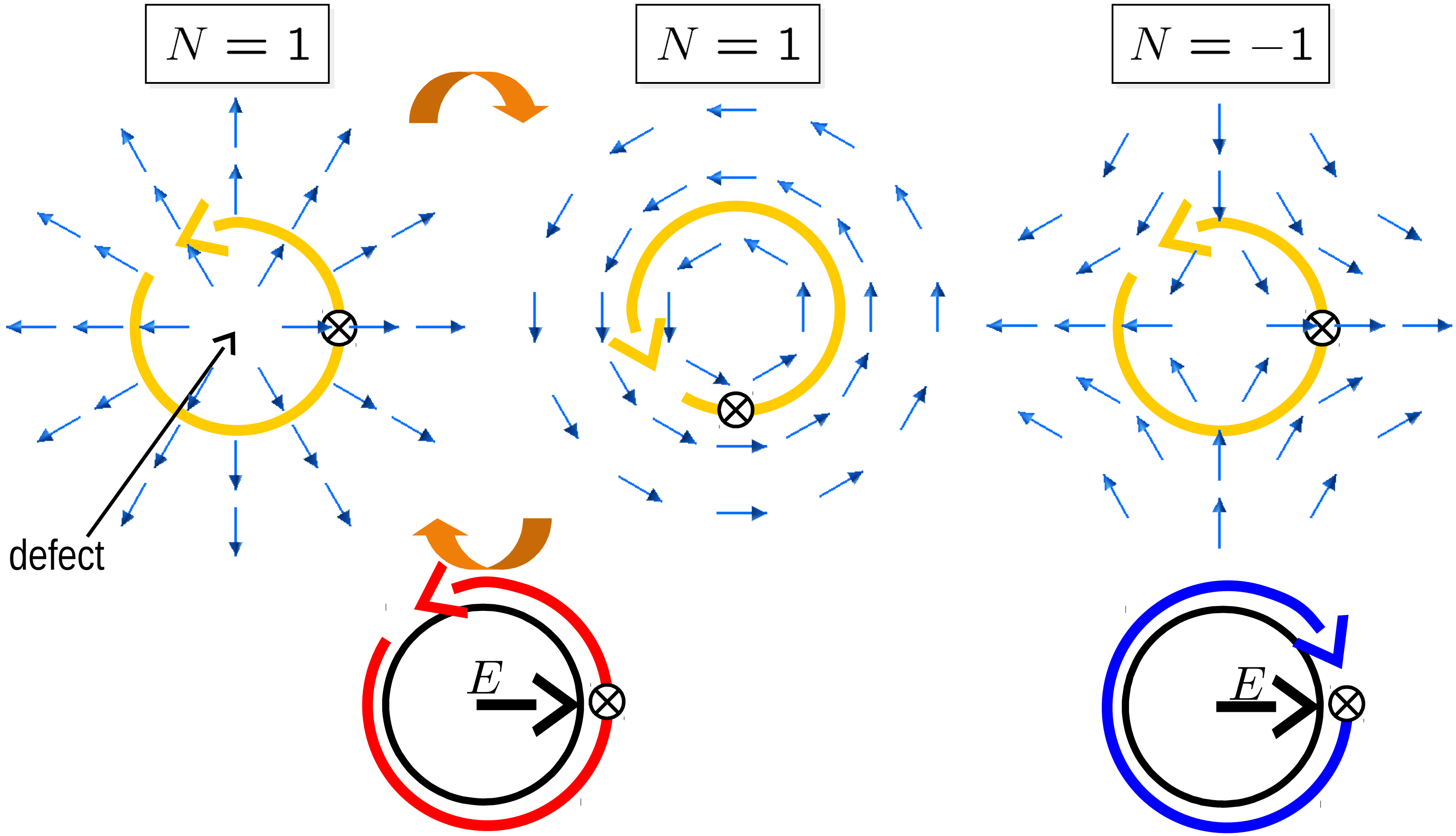
XY-spin system



In XY-spin system, local spin (order parameter) can be expressed by a point in a circle

$\longrightarrow$Order-parameter manifold

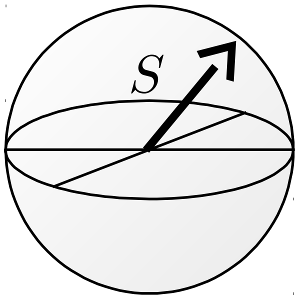
Topological defects and homotopy

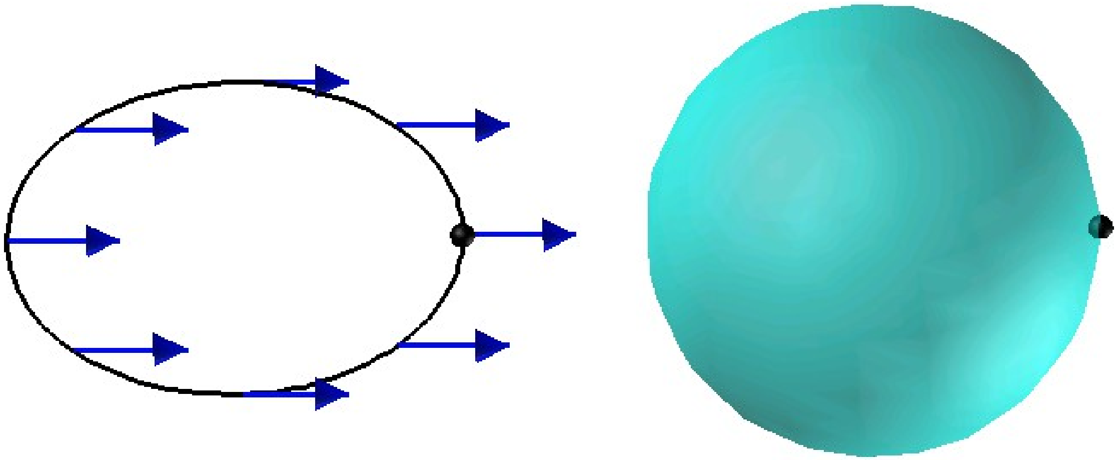


Topological excitations can be characterized by how many times the state rotates the circle along the closed path

Heisenberg-spin

Order parameter can be expressed by a point in a sphere





Topological excitations can never be stabilized

What is topological excitations?

Excitations in symmetry broken systems via phase transitions

Liquid → Solid transition (spontaneous symmetry breaking)

Liquid

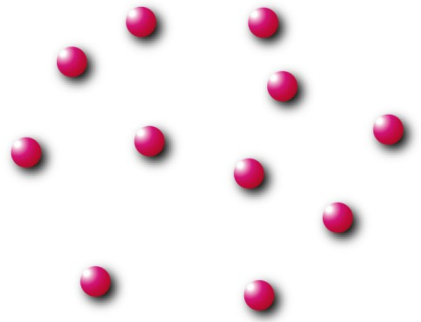
Solid (crystal)

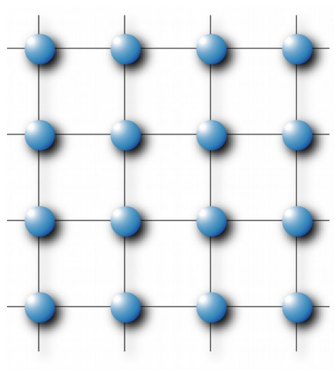
- Free energy is invariant under translational and rotational transformations
- System is also invariant under transformations

- Free energy is invariant under translational and rotational transformations
- System is not invariant under transformations (symmetry breaking)

What is topological excitations?

Liquid





Solid

Atoms are little influenced by other atoms.

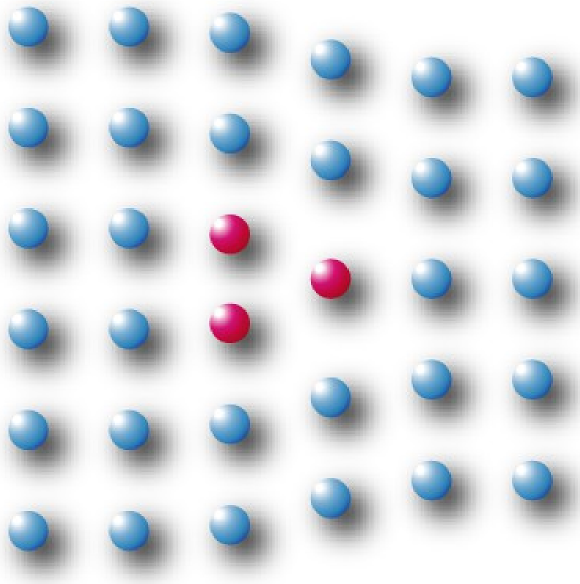
Positions and orientations of atoms are strongly affected by other atoms and fixed (spontaneous symmetry breaking).

What is topological excitations?

Topological excitations appear in symmetry broken systems

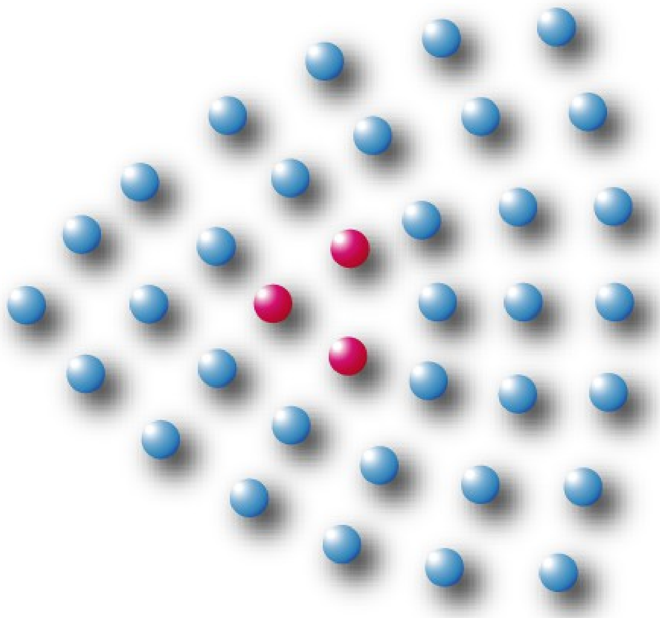
In crystal

dislocation



Topological excitation related to translational symmetry breaking

disclination



Topological excitation related to rotational symmetry breaking

$U(1)$ gauge symmetry breaking in BEC

Mean-field Hamiltonian at the zero temperature

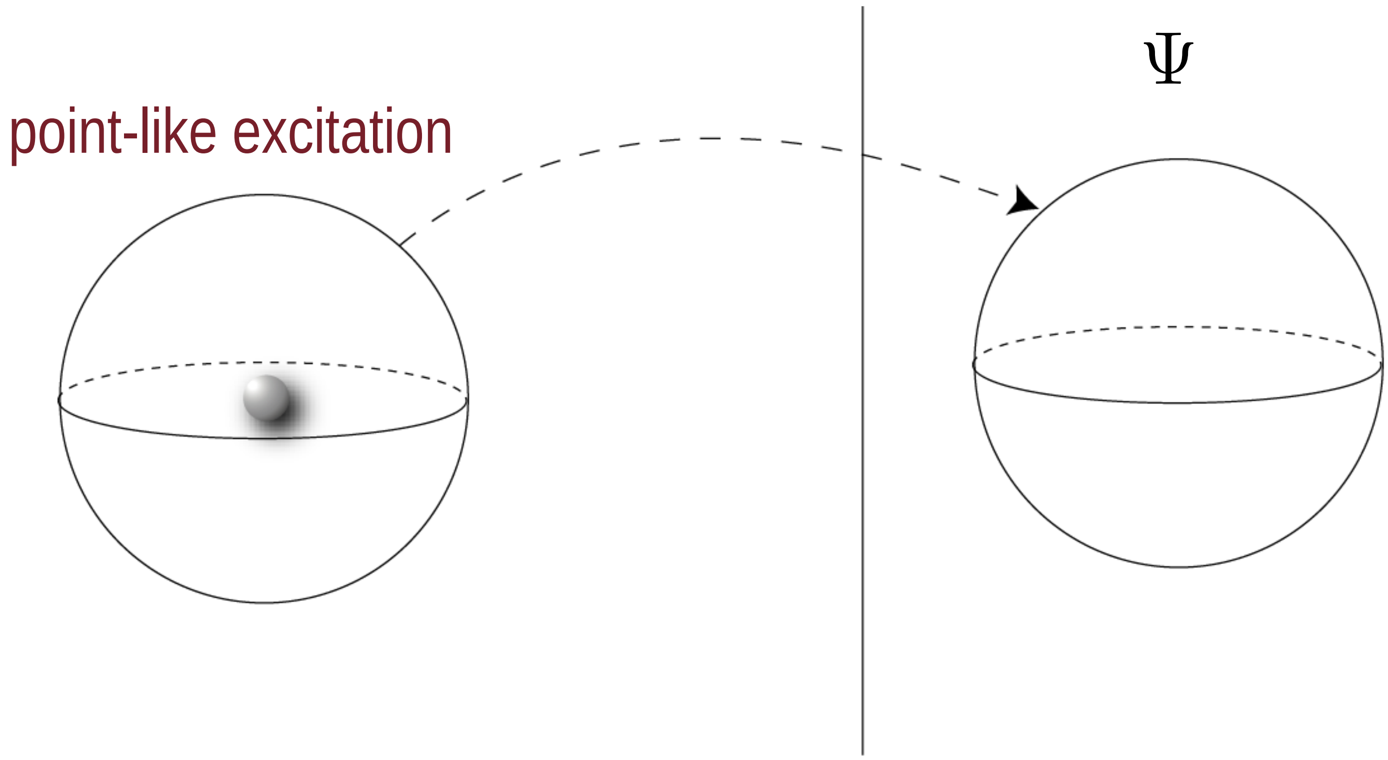
$$H = \int d\boldsymbol{x} \left[\frac{\hbar^2}{2M} \nabla \Psi^*(\boldsymbol{x}) \nabla \Psi(\boldsymbol{x}) + \frac{c_0}{2} |\Psi(\boldsymbol{x})|^4 \right]$$

$$\Psi(\boldsymbol{x}) = |\Psi(\boldsymbol{x})| \exp[i\varphi(\boldsymbol{x})]$$

$$\rho(\boldsymbol{x}) = |\Psi(\boldsymbol{x})|^2 \text{ : Fluid density}$$

$$\boldsymbol{v}(\boldsymbol{x}) = \frac{\hbar}{m} \nabla \varphi(\boldsymbol{x}) \text{ : Fluid velocity}$$

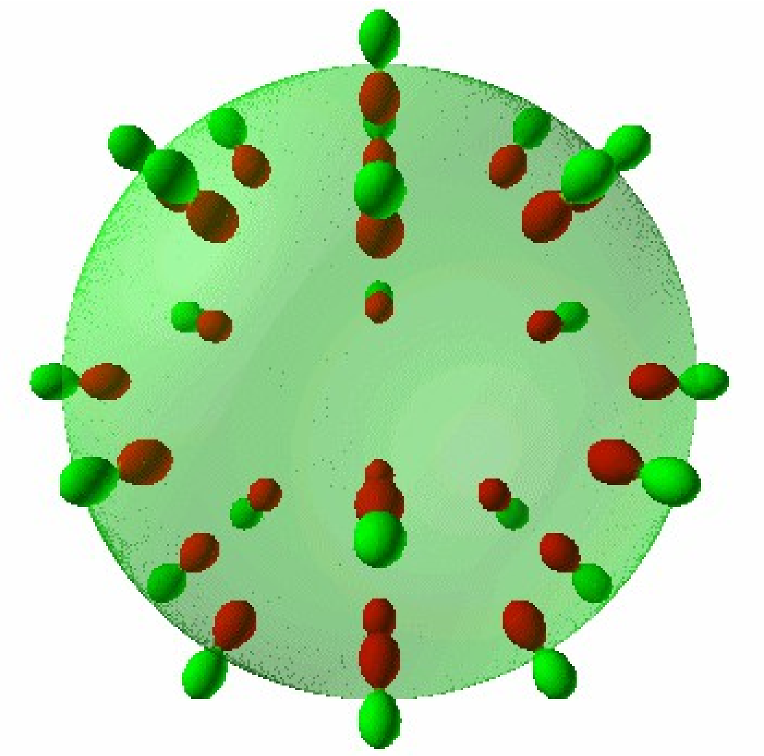
Point-like excitation



Point-like excitation

Polar phase

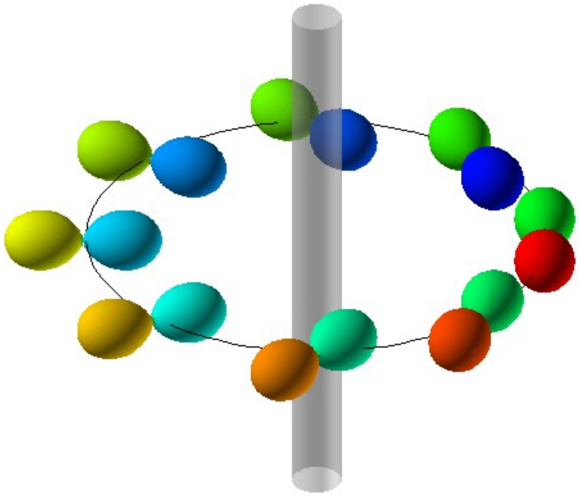
$$\frac{G}{H} \simeq \frac{U(1)_\varphi \times S_F^2}{(\mathbb{Z}_2)_{\varphi+F}}$$



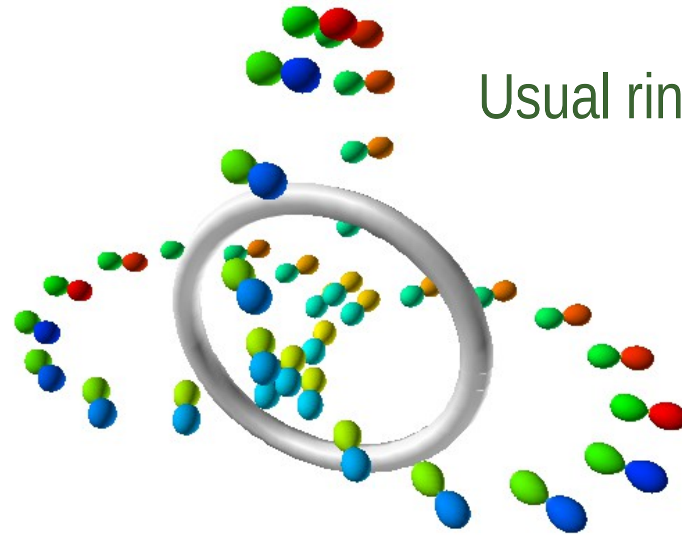
Point-like excitation cannot exist in Ferromagnetic phase

$$\frac{G}{H} \simeq SO(3)_{\varphi+F}$$

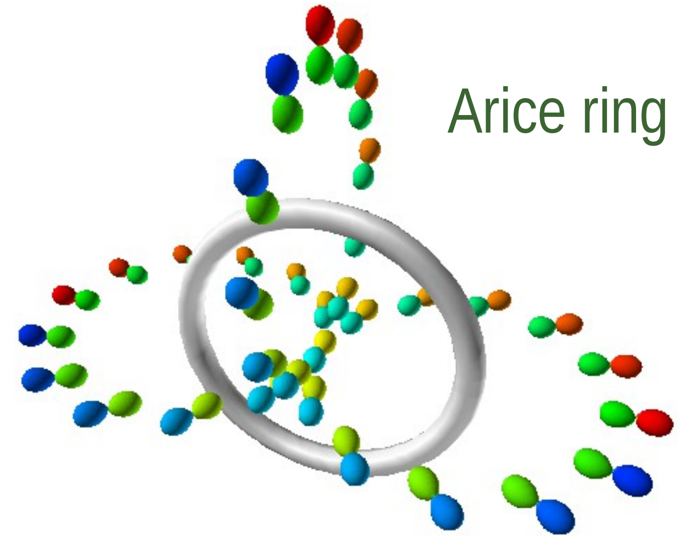
Vortex ring in polar phase



Vortex ring



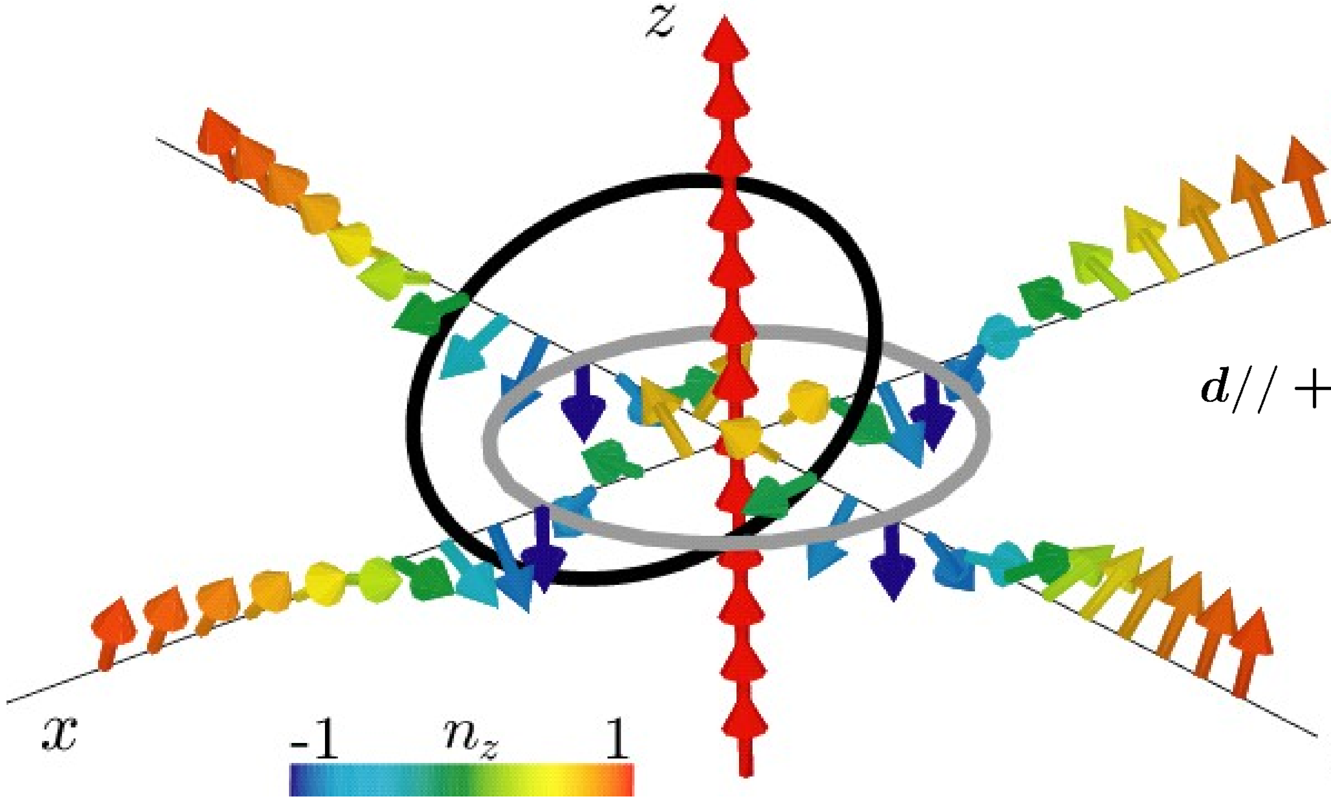
Usual ring

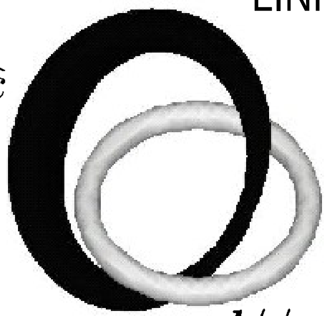


Arice ring

π_3 excitation and Hopf mapping

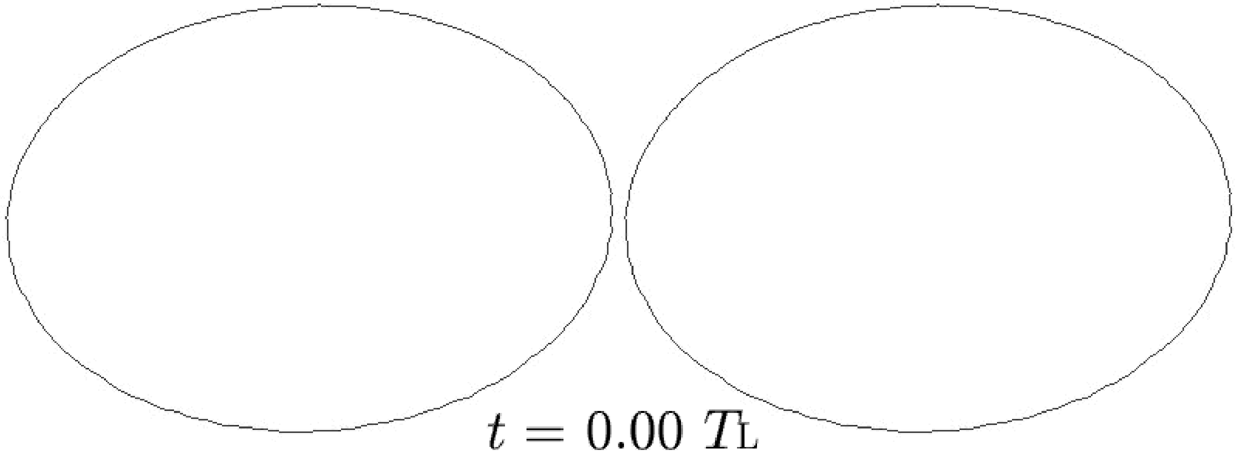
Y. Kawaguchi, et al., PRL **100**, 180403 (2008)

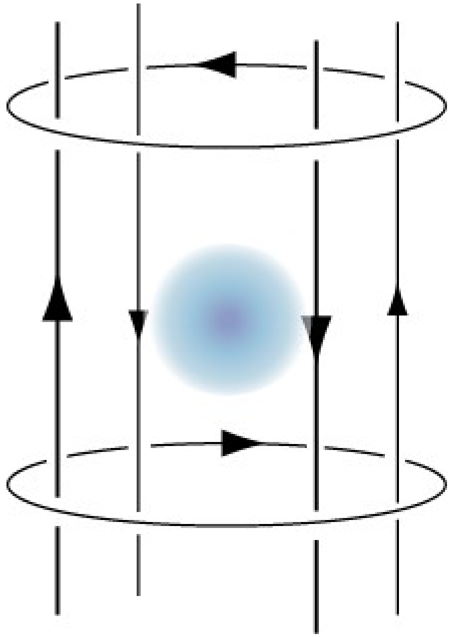


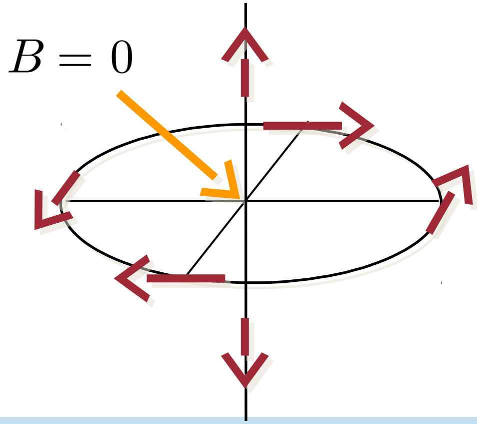


$$\pi_3 \left[\frac{U(1)_{\mathrm{G}} \times (S^2)_{\mathrm{S}}}{(\mathbb{Z}_2)_{\mathrm{G+S}}} \right] \cong (\mathbb{Z})_{\mathrm{S}}$$

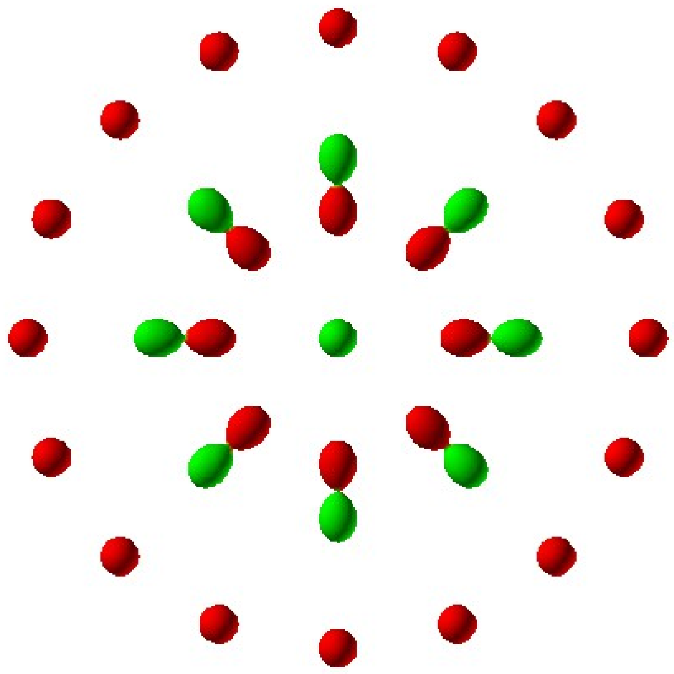
π_3 excitation and Hopf mapping

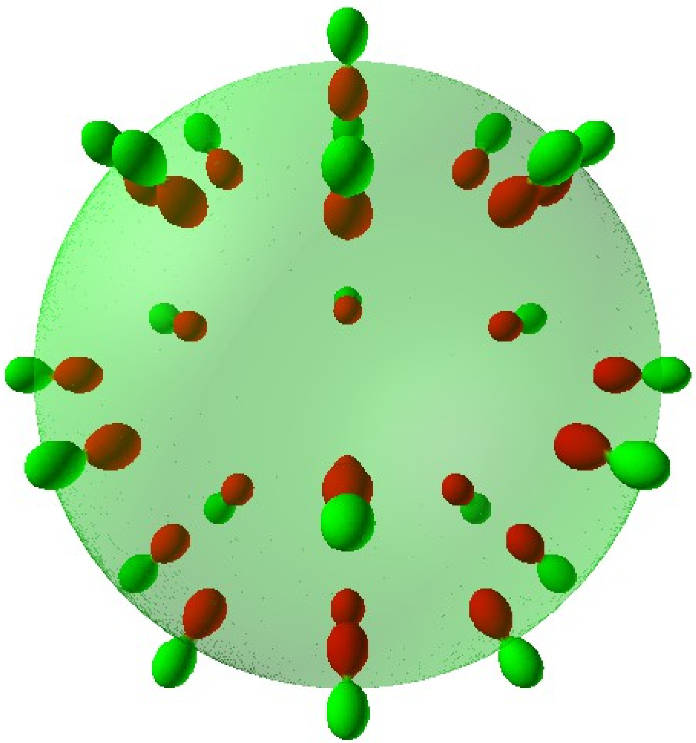






Point-like excitation and 2D skyrmion

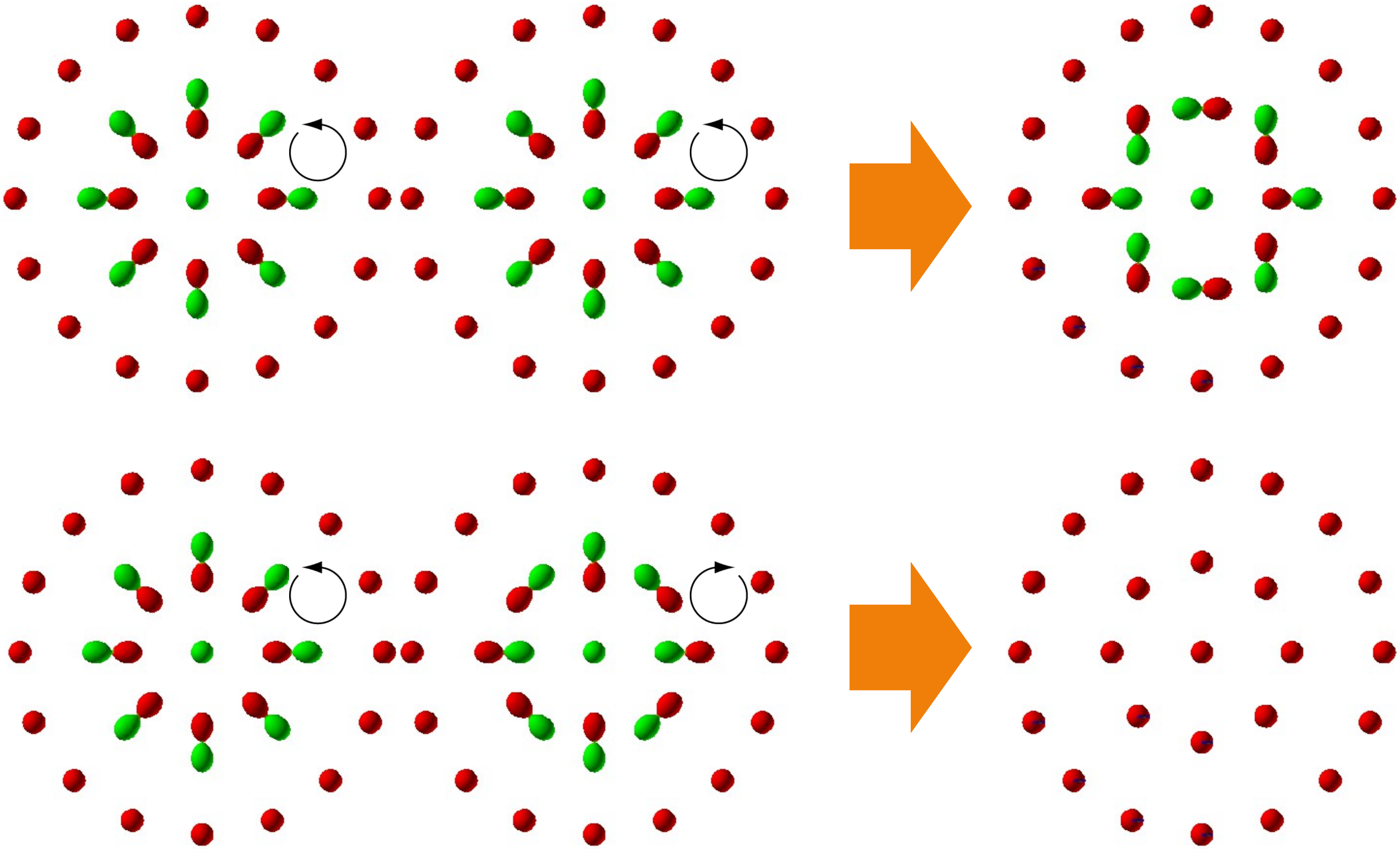


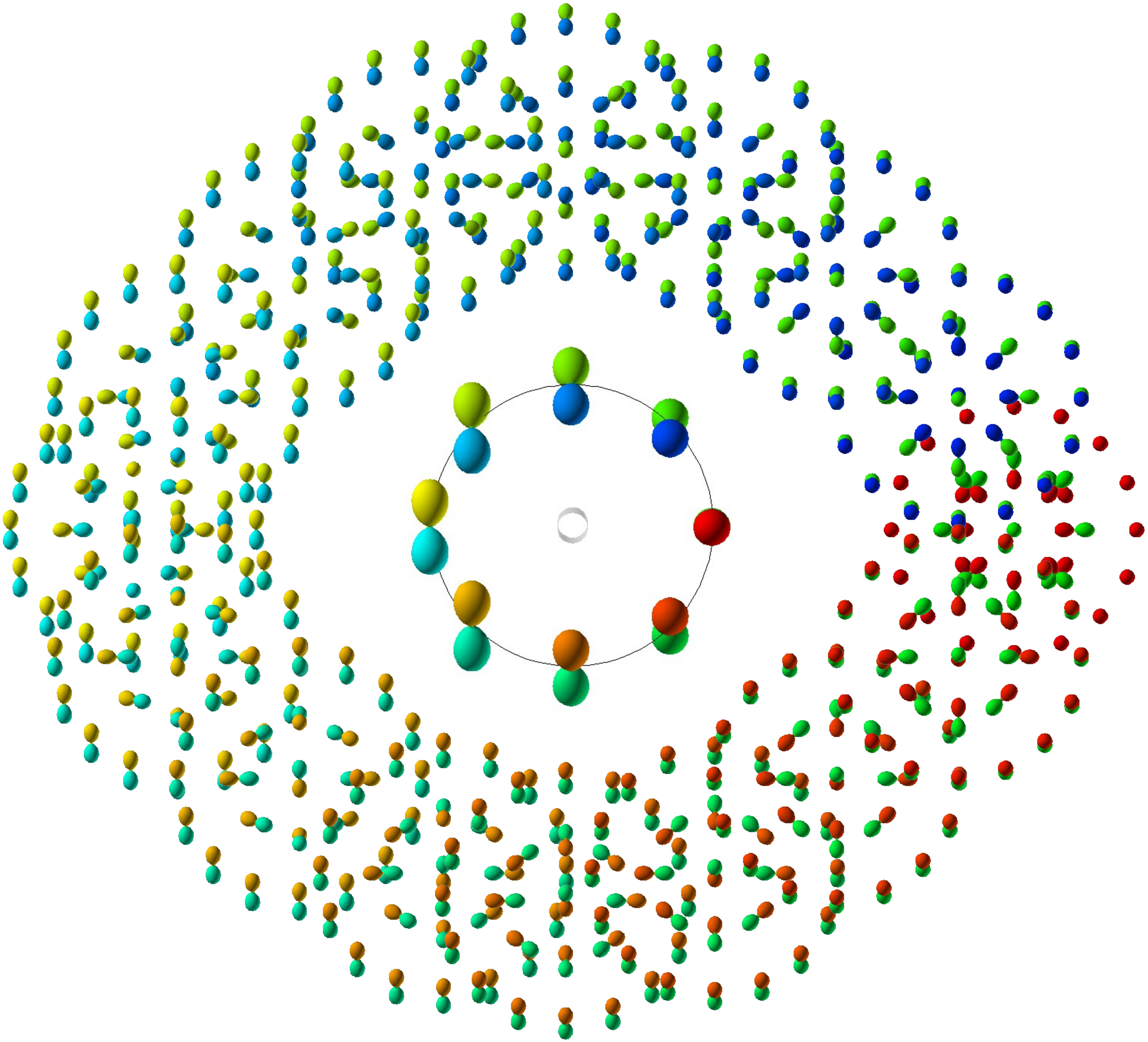


L. S. Leslie, et al. arXiv:0910.4918

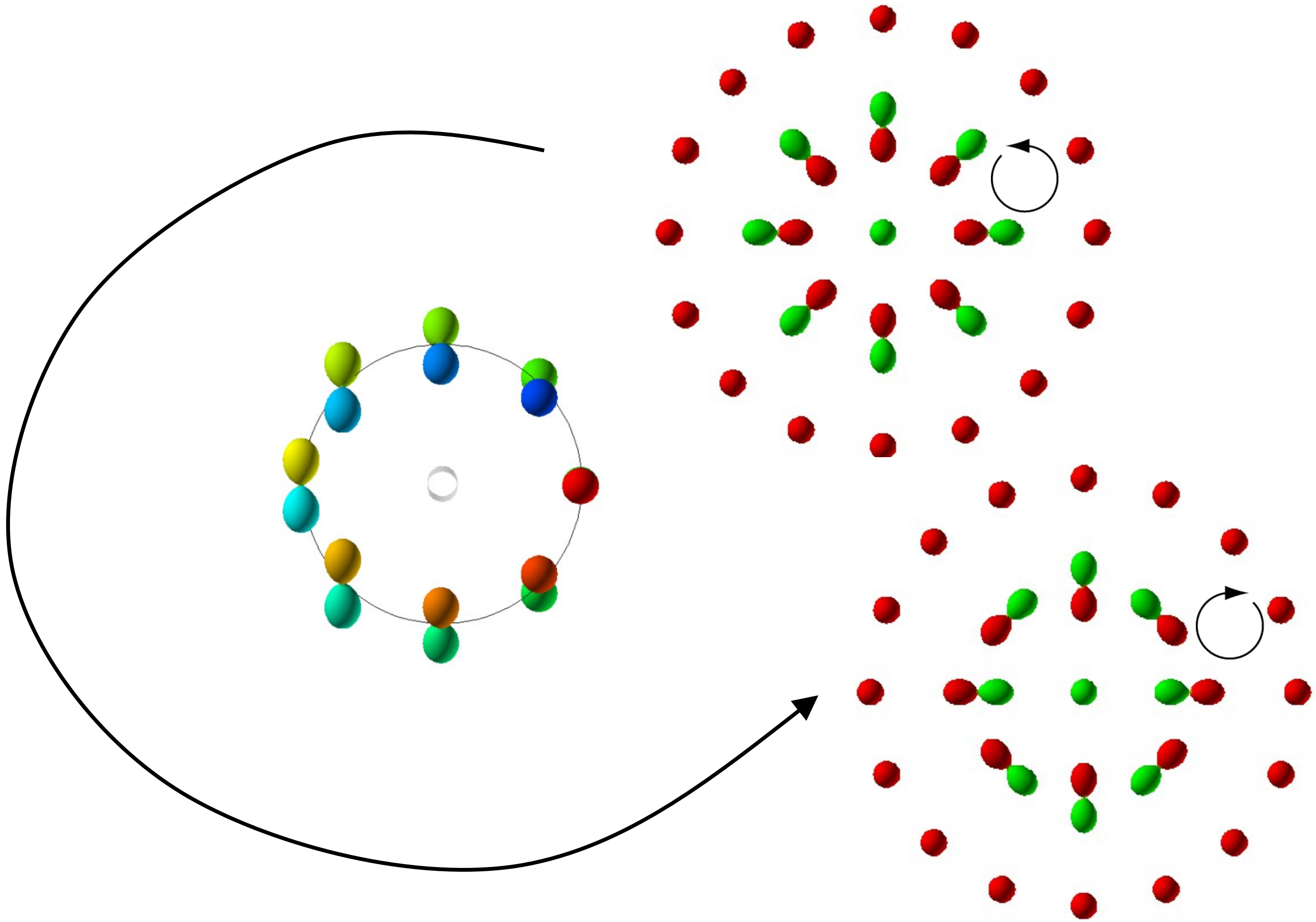
$$\pi_2\left[\frac{U(1)_{\mathrm{G}} \times (S^2)_{\mathrm{S}}}{(\mathbb{Z}_2)_{\mathrm{G+S}}}\right] \cong \pi_2[(S^2)_{\mathrm{S}}] \cong (\mathbb{Z})_{\mathrm{S}}$$

Two 2D skyrmion



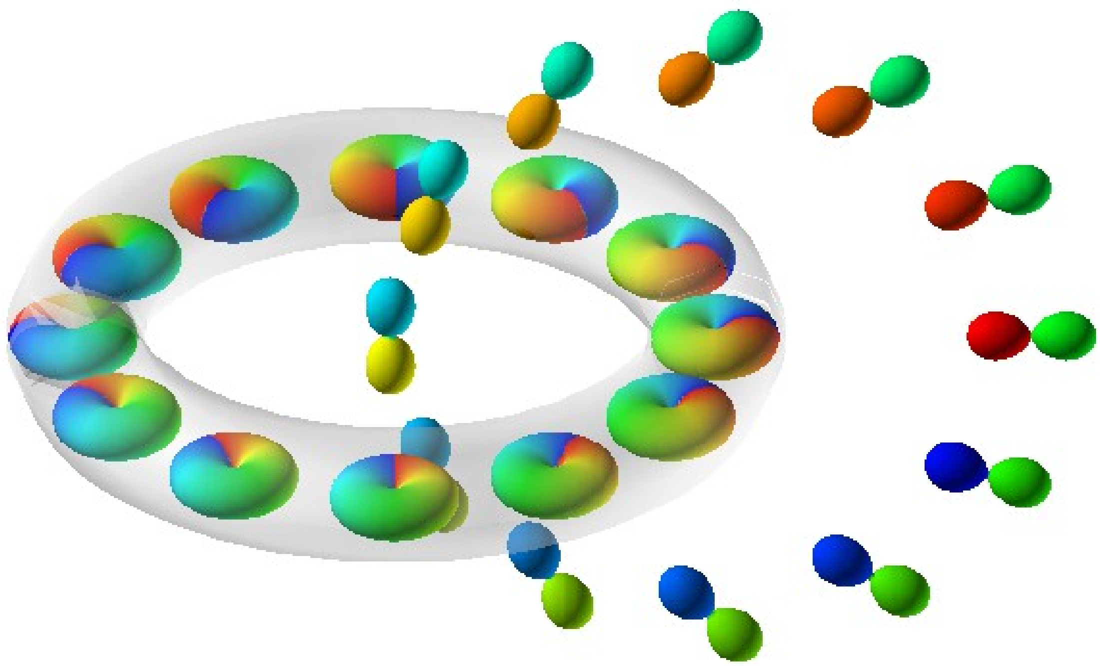


Inversion of topological invariant



Vorton excitation

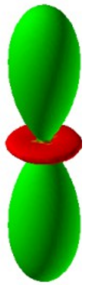
vorton



Spin-2 case

$$H = \int d\boldsymbol{x} \left[-\Psi_m^* \frac{\hbar^2}{2M} \nabla^2 \Psi_m + \frac{c_0}{2}\rho^2 + \frac{c_1}{2}\boldsymbol{F}^2 + \frac{c_2}{2}|A_{20}|^2 \right]$$

c_1



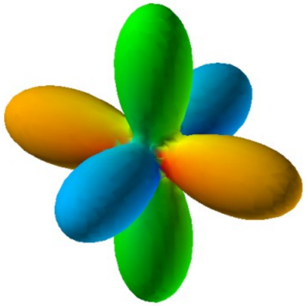
Uniaxial Nematic:

$$\Psi_{\mathrm{U}} = (0,0,1,0,0)^T$$

$$U(1)_\varphi \times \frac{S_F^2}{(\mathbb{Z}_2)_F}$$

Cyclic:

$$\Psi_{\mathrm{C}} = (1,0,0,\sqrt{2},1)^T/\sqrt{3}$$



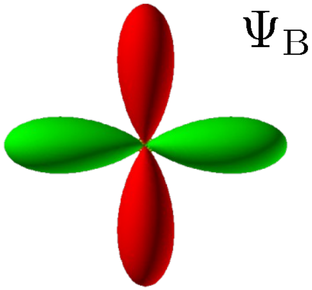
$$\frac{U(1)_\varphi \times SO(3)_F}{(T)_{\varphi+F}}$$

^{87}Rb

c_2

Biaxial Nematic:

$$\Psi_{\mathrm{B}} = (1,0,0,0,1)^T/\sqrt{2}$$

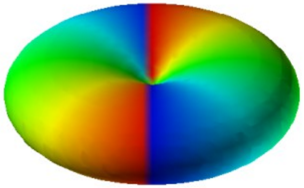


$$\frac{U(1)_\varphi \times SO(3)_F}{(D_4)_{\varphi+F}}$$

Ferromagnetic:

$$\Psi_{\mathrm{F}} = (1,0,0,0,0)^T$$

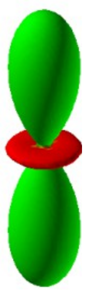
$$\frac{SO(3)_{\varphi+F}}{(\mathbb{Z}_2)_{\varphi+F}}$$



$$c_2 = 4c_1$$

Nematic phase of spin-2

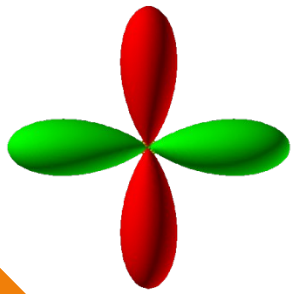
$$H = \int d\boldsymbol{x} \left[-\Psi_m^* \frac{\hbar^2}{2M} \nabla^2 \Psi_m + \frac{c_0}{2}\rho^2 + \frac{c_1}{2}\boldsymbol{F}^2 + \frac{c_2}{2}|A_{20}|^2 \right]$$



Uniaxial Nematic:

$$\Psi_{\mathrm{U}} = (0,0,1,0,0)^T$$

$$U(1)_\varphi \times \frac{S_F^2}{(\mathbb{Z}_2)_F}$$

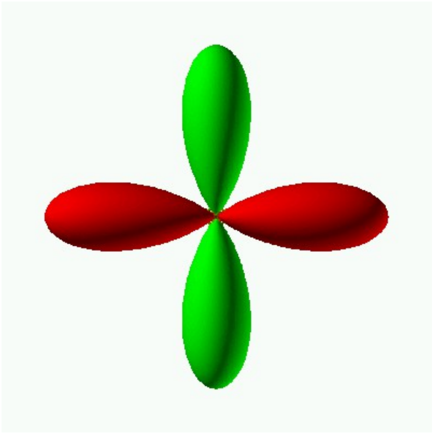


Biaxial Nematic:

$$\Psi_{\mathrm{B}} = (1,0,0,0,1)^T/\sqrt{2}$$

$$\frac{U(1)_\varphi \times SO(3)_F}{(D_4)_{\varphi+F}}$$

Two states are degenerate via another continuous degree of freedom

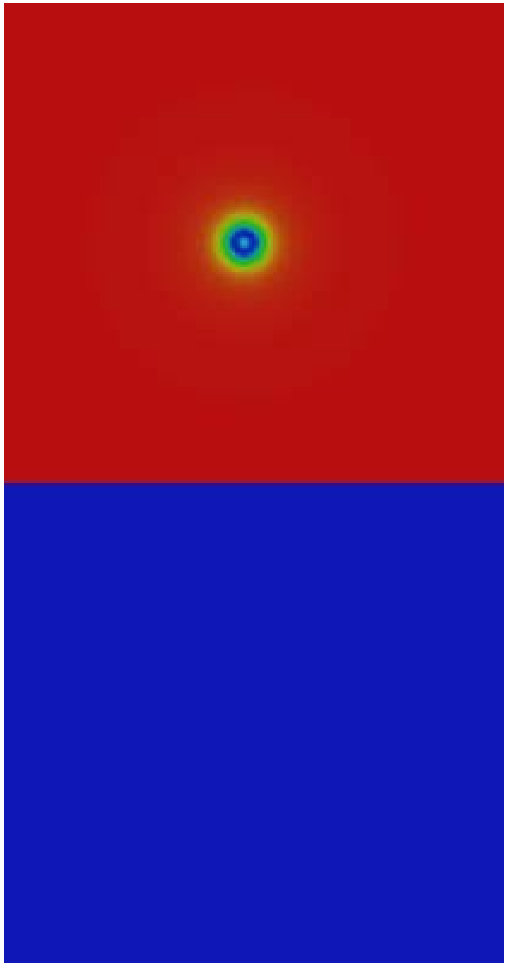


New order-parameter manifold

$$\frac{U(1)_\varphi \times S_F^4}{(\mathbb{Z}_2)_{\varphi+F}}$$

Quasi-Nambu-Goldstone current

S. Uchino, et al., PRL in press



Decay of vortex in biaxial nematic phase

Emission of quasi-Nambu-Goldstone current

Cyclic State vs. Singlet-trio Condensed State

For $c_1 > 0,\ c_2 > 0$

M. Koashi, and M. Ueda. PRL **84**, 1066 (2000)

Singlet-trio condensed state (only $\mathrm{U}(1)$ is broken)

$$|\Psi\rangle = \left[e^{i\varphi} \left(\frac{\sqrt{2}\hat{a}_0(\hat{a}_0^{\dagger 2} - 3a_1^\dagger a_{-1}^\dagger - 6a_2^\dagger a_{-2}^\dagger) + 3\sqrt{3}(a_1^{\dagger 2} a_{-2}^\dagger + a_{-1}^{\dagger 2} a_2^\dagger)}{\sqrt{35}} \right) \right]^{N/3} |0\rangle$$

Transition occurs under ~1μG

Cyclic state ($\mathrm{U}(1) \times \mathrm{SO}(3)$ is broken)

$$|\Psi\rangle = \left[\sum_m \Psi_m a_m^\dagger \right]^N |0\rangle$$

$$\Psi = e^{i\varphi} e^{-i\hat{\boldsymbol{F}}\cdot\boldsymbol{\alpha}} \begin{pmatrix} i/2 \\ 0 \\ 1/\sqrt{2} \\ 0 \\ i/2 \end{pmatrix}$$

Nematic State vs. Singlet-pair Condensed State

For $c_1 > 0,\ c_2 < 0$

M. Koashi, and M. Ueda. PRL **84**, 1066 (2000)

Singlet-pair condensed state (only $\mathrm{U}(1)$ is broken)

$$|\Psi\rangle = \left[e^{i\varphi} \left(\frac{\hat{a}_0^{\dagger 2} - 2a_1^\dagger a_{-1}^\dagger + a_2^\dagger a_{-2}^\dagger}{\sqrt{5}} \right) \right]^{N/2} |0\rangle$$

Transition occurs under ~1μG

Nematic state $(\mathrm{U}(1) \times \mathrm{SO}(3)$ is broken)

$$|\Psi\rangle = \left[\sum_m \Psi_m a_m^\dagger \right]^N |0\rangle \qquad \Psi = e^{i\varphi} e^{-i\hat{\boldsymbol{F}}\cdot\boldsymbol{\alpha}} \begin{pmatrix} 0 \\ 0 \\ 1 \\ 0 \\ 0 \end{pmatrix} \text{ or } \begin{pmatrix} 1 \\ 0 \\ 0 \\ 0 \\ 1 \end{pmatrix} / \sqrt{2}$$